



# Generation and Characterization of Structured Partially Coherent Light

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GIQSUL – Department of Physics  
Federal University of Santa Catarina  
Florianópolis

February, 2025



# WORLD MAP



Map is not to Scale  
infoandopinion.com

Antarctica  
W 0° E

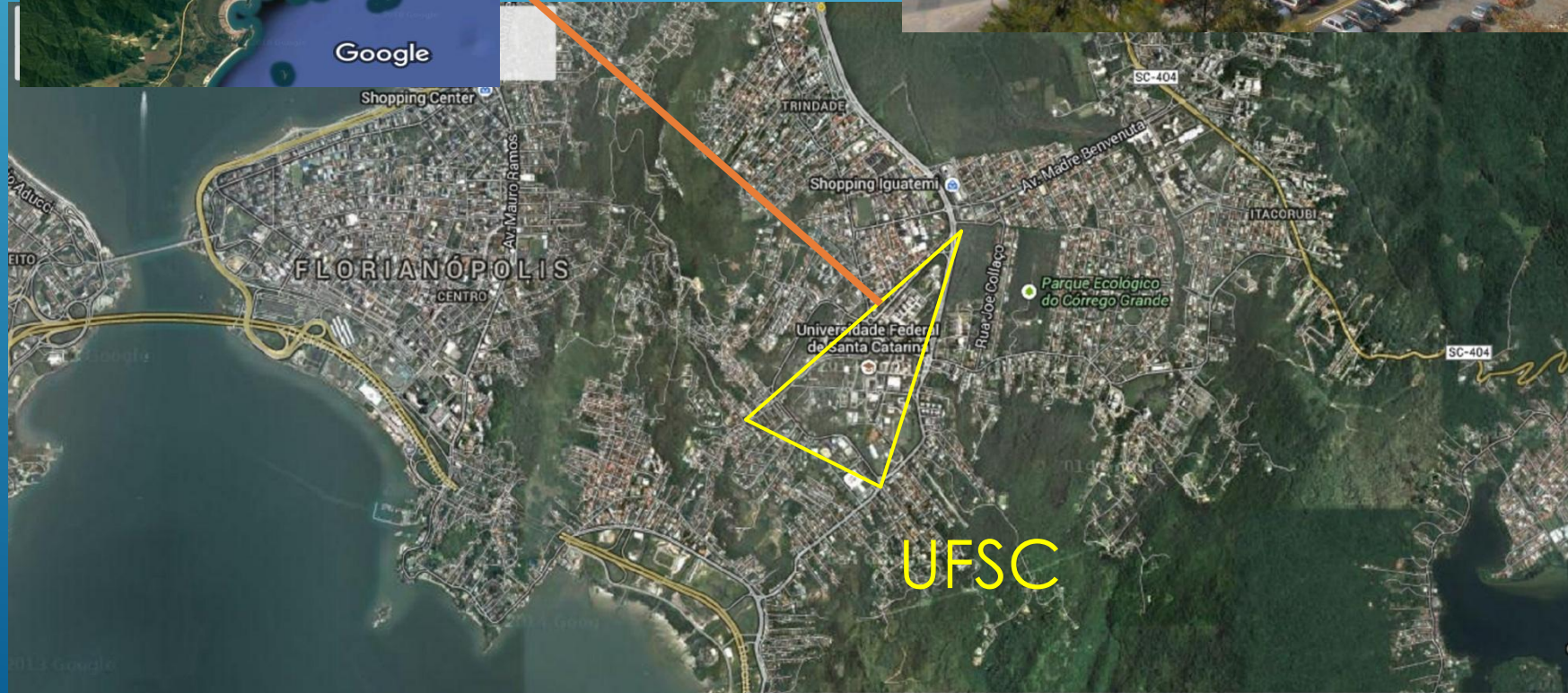
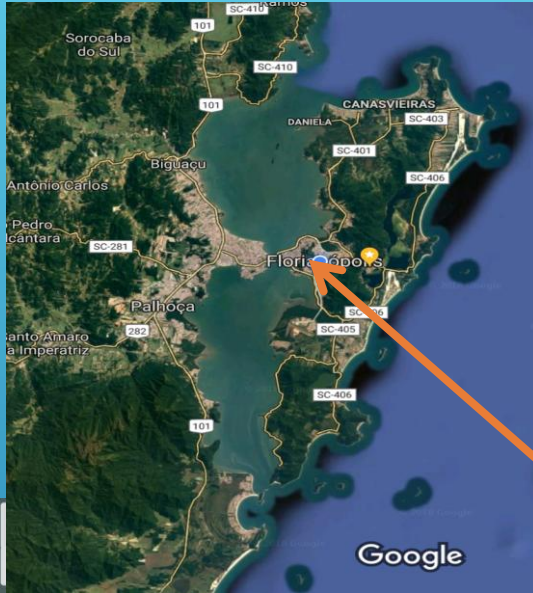


Santa Catarina

# Florianópolis, SC

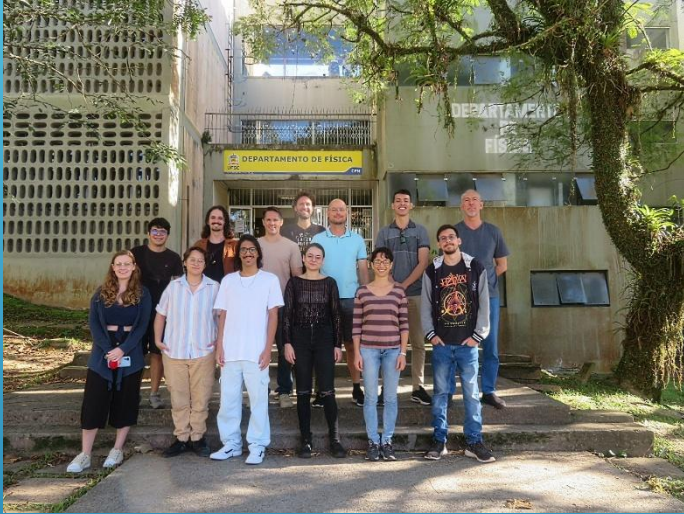


# Federal University of Santa Catarina UFSC

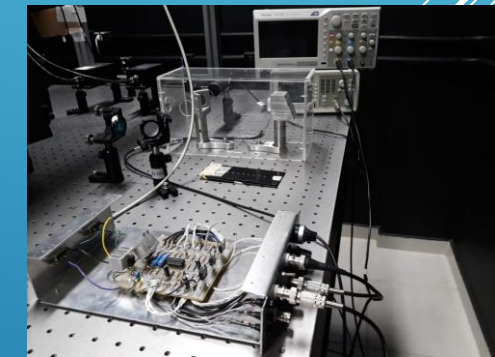
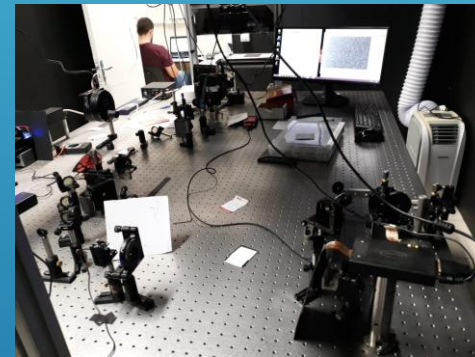
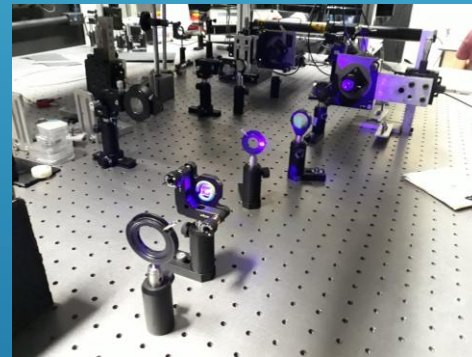
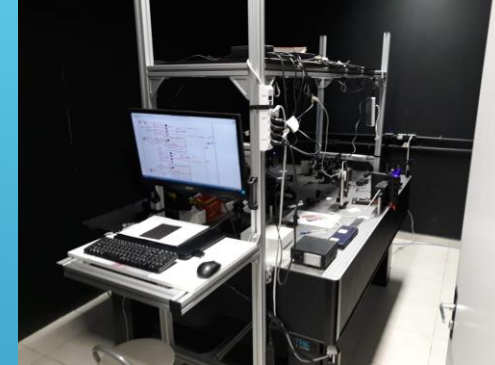
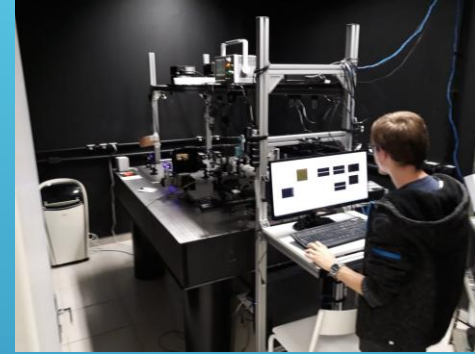


# GIQSUL at UFSC

## Group of Quantum Information



## Quantum Optics Laboratory




## Optical Processing Laboratory



# GIQSUL Research

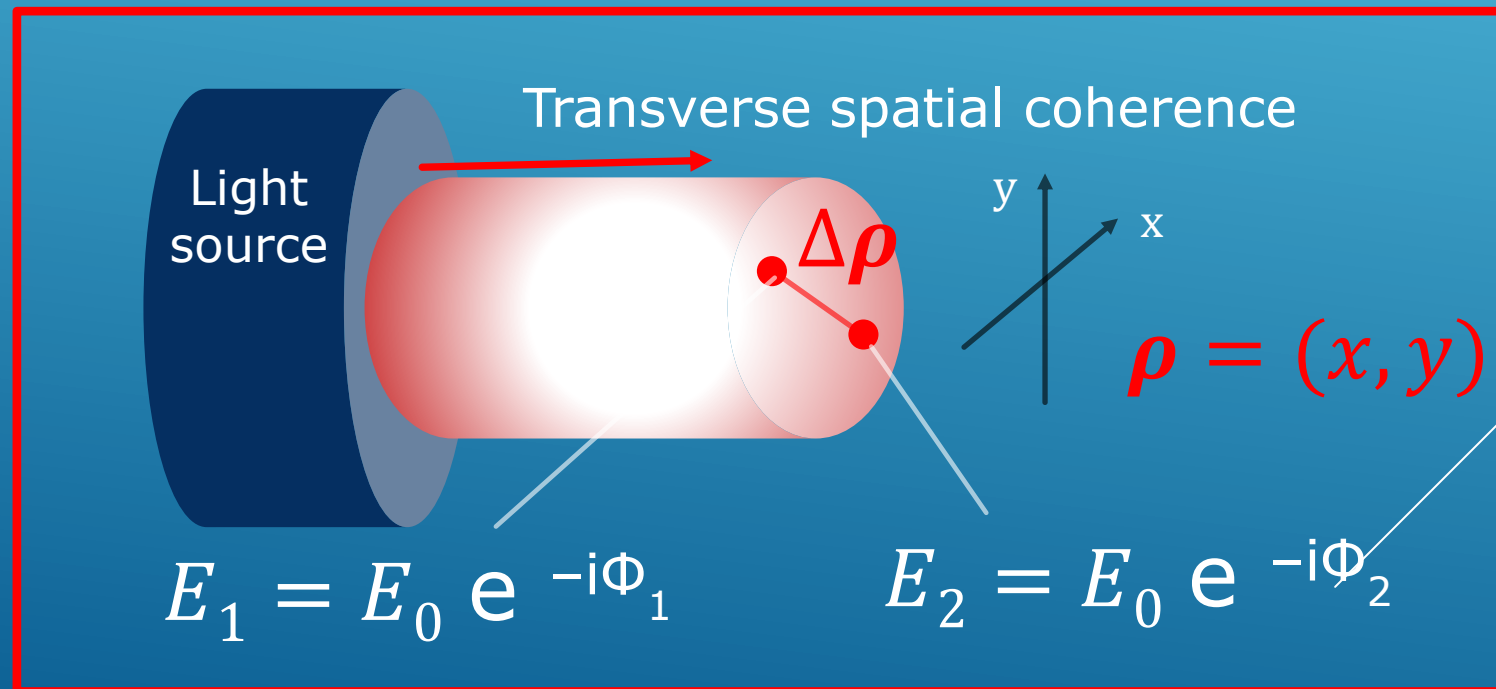
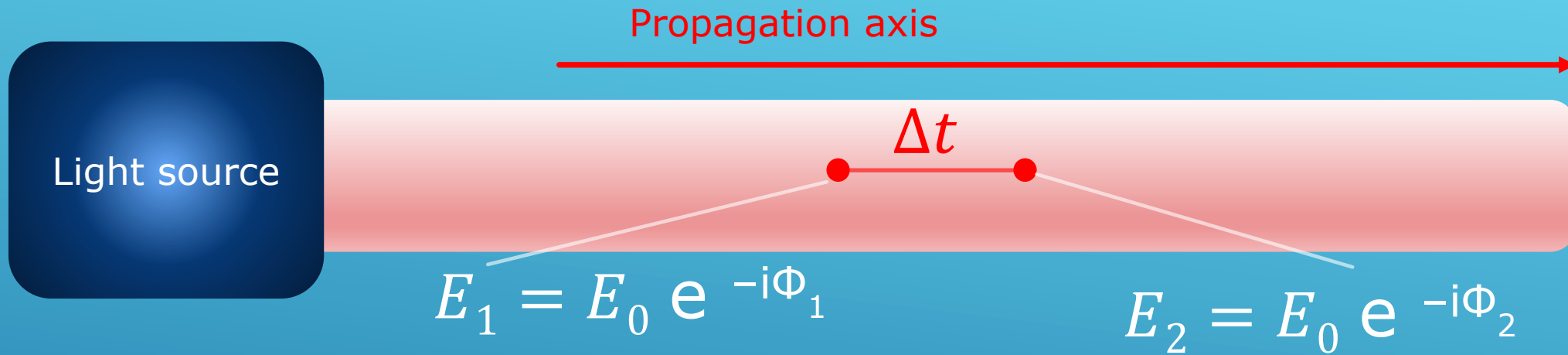
- **Transverse Spatial Entanglement with Parametric Downconversion**  
Spatial correlations in parametric down-conversion, Walborn et al., Physics Reports 495, 87-139 (2010)
- **Quantum Computing Theory (Duzzioni team)**  
Quantum computation in continuous time using dynamic invariants, Sarandi et al., Physics Letters A, 375, 3343-3347 (2011)
- **Nonlinear Optics**  
Conservation of orbital angular momentum in stimulated downconversion, Caetano et al. Phys. Rev. A 66, 041801(R) (2002)
- **Optical Processing**  
An optical processor for matrix-by-vector multiplication: an application to the distance geometry problem in 1D, Hengeveld et al. Journal of Optics 24, 015701 (2021)

# Outline

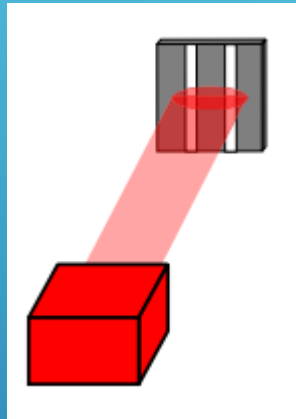
- Introduction to optical coherence
  - The Gaussian Schell Model Beam (GSM)
  - The Twisted Gaussian Schell Model Beam (TGSM)
  - Motivation for TGSM beams
  - Generation of TGSM beams
  - Quantum effects with TGSM beams
  - StimPDC with TGSM beams
  - Conclusions and perspectives
- 



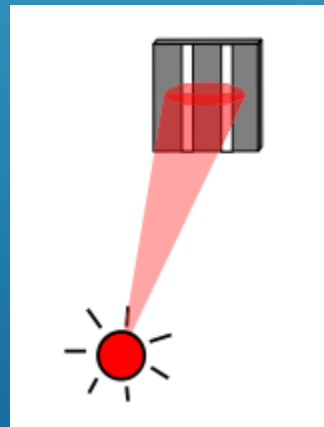
# Temporal and transverse spatial coherence



# Measuring transverse spatial coherence with double-slit interference



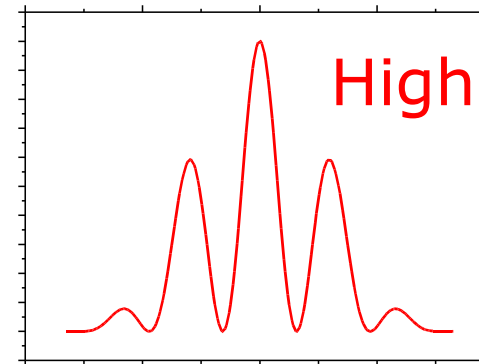
Laser



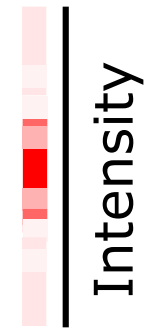
Extended Source



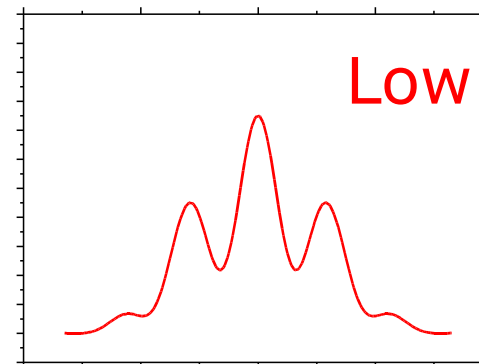
Intensity



Position



Intensity

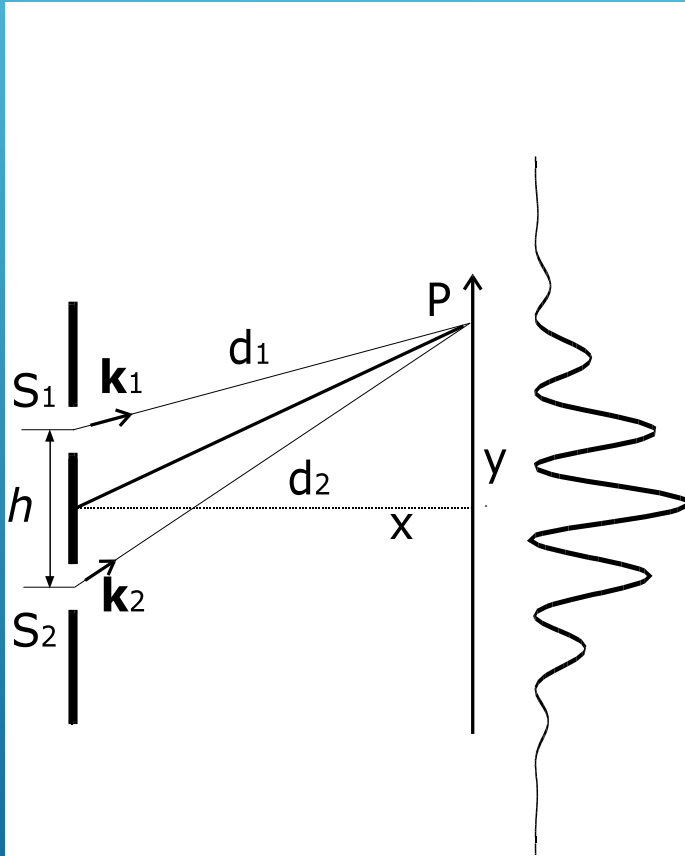


Position

Contrast or visibility  $\Rightarrow$   $|\mu_{12}| = (I_{\max} - I_{\min}) / (I_{\max} + I_{\min})$

# Coherence and Double-slit interference

Almost monochromatic light  $\sim$  single frequency



Intensity distribution

$$I(p) = I_0(y)(1 + |\mu_{12}| \cos[k(d_2 - d_1) + \varphi])$$

$$k = 2\pi/\lambda$$

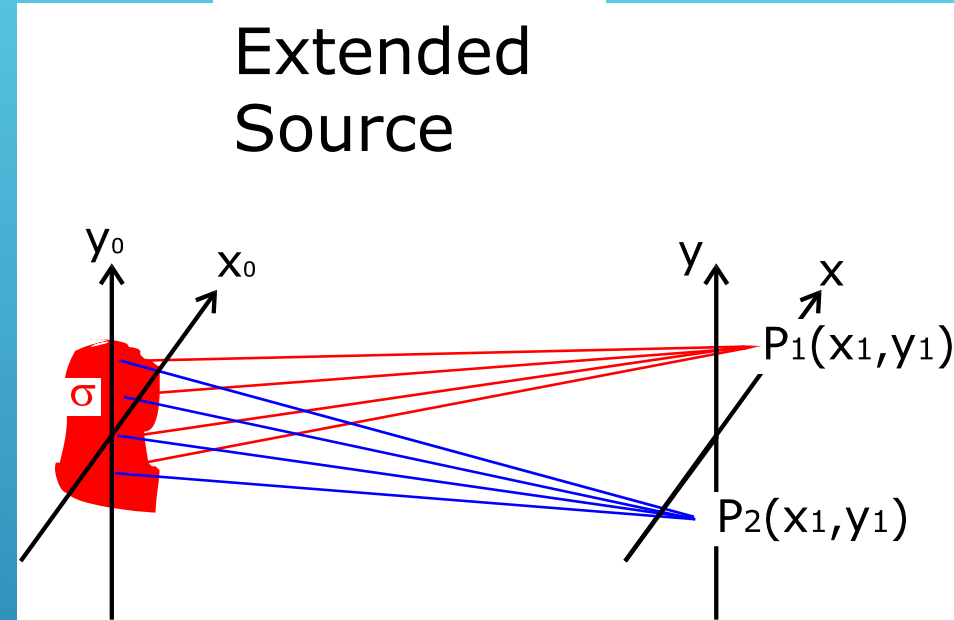
$$d_1 = \sqrt{x^2 + \left(y - \frac{h}{2}\right)^2} ; d_2 = \sqrt{x^2 + \left(y + \frac{h}{2}\right)^2}$$

$\mu_{12}$  Normalized degree of mutual coherence

# Van Cittert-Zernike theorem

Cross spectral density

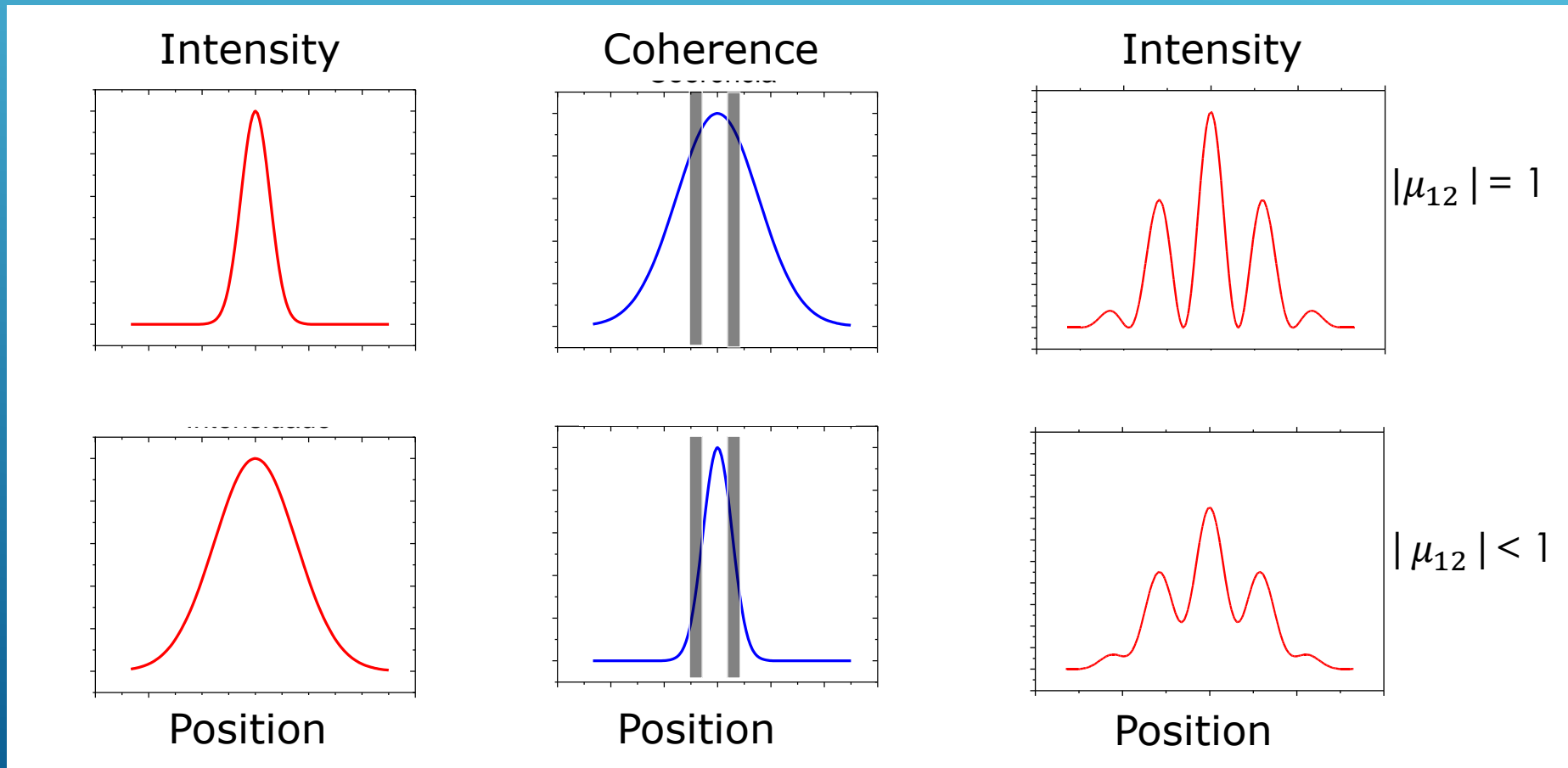
$$\mu_{12} = \frac{\langle E_1(x_1, y_1) E_2^*(x_2, y_2) \rangle}{\int_{\sigma} dx_0 dy_0 I(x_0, y_0)}$$



$$\mu_{12}[(x_2 - x_1), (y_2 - y_1)] = \frac{e^{i\alpha_{12}} \int_{\sigma} dx_0 dy_0 I(x_0, y_0) e^{i\frac{k}{R}[x_0(x_2 - x_1) + y_0(y_2 - y_1)]}}{\int_{\sigma} dx_0 dy_0 I(x_0, y_0)}$$

# Van Cittert-Zernike theorem

$$\mu_{12}[(x_2 - x_1), (y_2 - y_1)] = \frac{e^{i\alpha_{12}} \int_{\sigma} dx_0 dy_0 I(x_0, y_0) e^{i\frac{k}{R}[x_0(x_2 - x_1) + y_0(y_2 - y_1)]}}{\int_{\sigma} dx_0 dy_0 I(x_0, y_0)}$$



# The Gaussian Schell Model (GSM)

## THE MULTIPLE PLATE ANTENNA

by

ALLAN CARTER SCHELL

S.B., Massachusetts Institute of Technology  
(1956)

S.M., Massachusetts Institute of Technology  
(1956)

SUBMITTED IN PARTIAL FULFILLMENT OF THE  
REQUIREMENTS FOR THE DEGREE OF  
DOCTOR OF SCIENCE

at the

MASSACHUSETTS INSTITUTE OF TECHNOLOGY  
September, 1961

## THE MULTIPLE PLATE ANTENNA

by

ALLAN CARTER SCHELL

Submitted to the Department of Electrical Engineering on August 21, 1961  
in partial fulfillment of the requirements for the degree of Doctor of Science

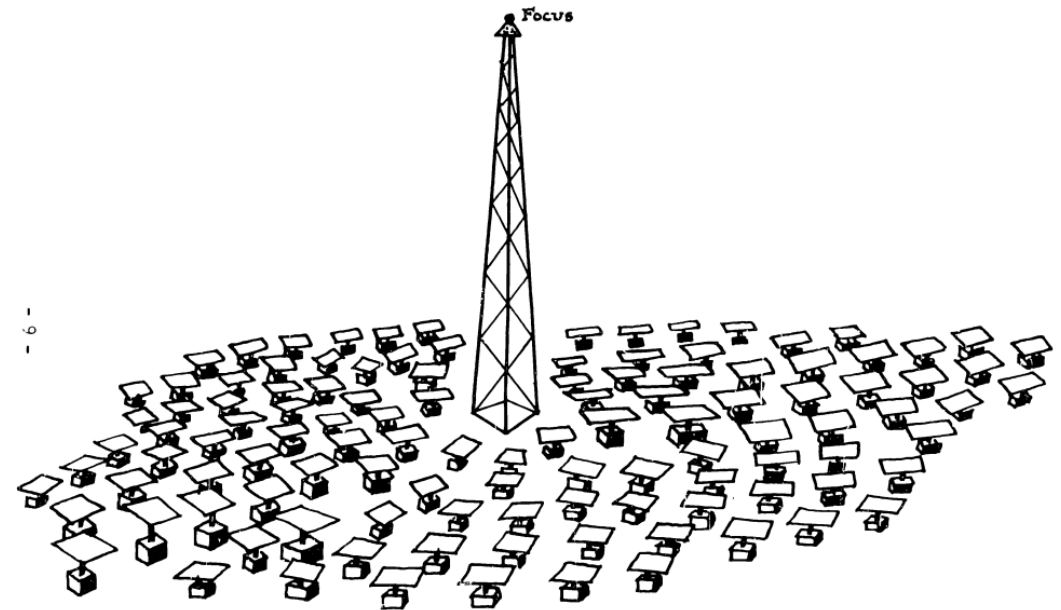


Figure 1

*A Sketch of a Multiple Plate Radio Astronomy Antenna*

# The Gaussian Schell Model (GSM) beams

GAUSSIAN SCHELL-MODEL BEAMS

Leonard Mandel · Emil Wolf *Editors*

Coherence and  
Quantum Optics V

Ronald J. Sudol

Department of Physics and Astronomy  
University of Rochester, Rochester, N.Y. 14627, USA



Ari T. Friberg

Department of Technical Physics  
Helsinki University of Technology  
SF-02150 Espoo 15, Finland

Proceedings of the Fifth Rochester Conference on  
Coherence and Quantum Optics held at the  
University of Rochester, June 13-15, 1983

# The Gaussian Schell Model (GSM) beams

In the Schell-model approximation the source cross-spectral density function takes the form

$$W(\underline{\rho}_1, 0; \underline{\rho}_2, 0) = [I(\underline{\rho}_1, 0)I(\underline{\rho}_2, 0)]^{\frac{1}{2}} \mu(\underline{\rho}_1 - \underline{\rho}_2; 0)$$

$$I(\underline{\rho}, 0) = A \exp\{-\rho^2/2\sigma_I^2\}$$

$$\mu(\underline{\rho}'; 0) = \exp\{-\rho'^2/2\sigma_\mu^2\}$$

GAUSSIAN SCHELL-MODEL BEAMS

Ronald J. Sudol

Department of Physics and Astronomy  
University of Rochester, Rochester, N.Y. 14627, USA

Ari T. Friberg

Department of Technical Physics  
Helsinki University of Technology  
SF-02150 Espoo 15, Finland

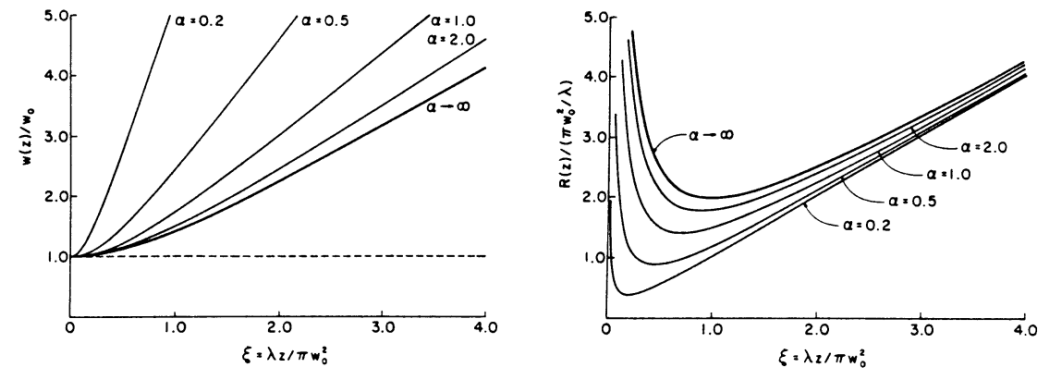


Fig. 1. The behavior of the beam radius  $w(z)$  and the radius of curvature  $R(z)$  of a Gaussian Schell-model beam as a function of the dimensionless variable  $\xi = \lambda\pi / \pi w_0^2$  for several values of the parameter  $\alpha = \sigma_\mu^2 / 2\sigma_I^2$ . The region  $\alpha \ll 1$  correspond to globally incoherent beams (Gaussian quasi-homogeneous beams), whereas the limit  $\alpha \rightarrow \infty$  represents a fully coherent Gaussian laser beam.



# The Gaussian Schell Model (GSM) beams

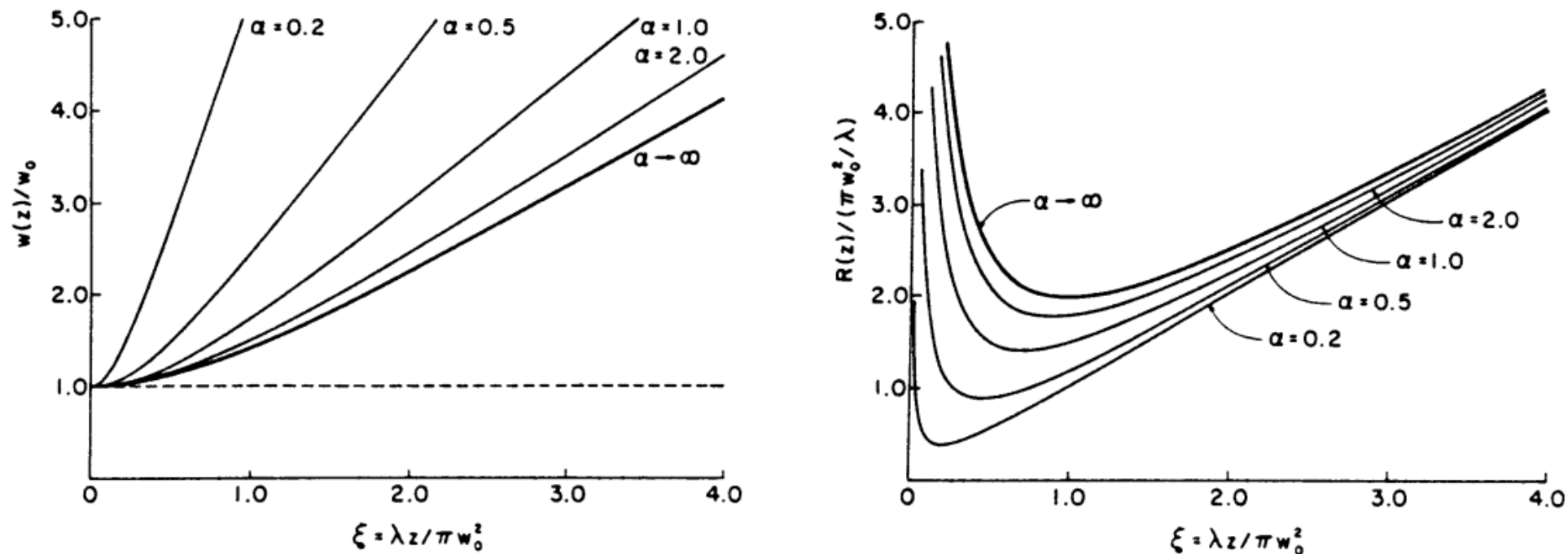


Fig. 1. The behavior of the beam radius  $w(z)$  and the radius of curvature  $R(z)$  of a Gaussian Schell-model beam as a function of the dimensionless variable  $\xi = \lambda \pi / \pi w_0^2$  for several values of the parameter  $\alpha = \sigma_{\mu} / 2\sigma_I$ . The region  $\alpha \ll 1$  correspond to globally incoherent beams (Gaussian quasi-homogeneous beams), whereas the limit  $\alpha \rightarrow \infty$  represents a fully coherent Gaussian laser beam.

# The Twisted Gaussian Schell Model (TGSM) beams

R. Simon and N. Mukunda

Vol. 10, No. 1/January 1993/J. Opt. Soc. Am. A 95

## Twisted Gaussian Schell-model beams



**Rajiah Simon**



**Narasimhaiengar Mukunda**

# The Twisted Gaussian Schell Model (TGSM) beams

R. Simon and N. Mukunda

Vol. 10, No. 1/January 1993/J. Opt. Soc. Am. A

95

## Twisted Gaussian Schell-model beams

We may ask, What is the most general Gaussian cross-spectral density (in a transverse plane) that is invariant under arbitrary rotations about the  $z$  axis? The answer

$$\begin{aligned} E(\boldsymbol{\rho}_1, \boldsymbol{\rho}_2) &= a_1(\rho_1^2 + \rho_2^2) + \gamma \boldsymbol{\rho}_1 \cdot \boldsymbol{\rho}_2 \\ &\quad + ia_2(\rho_1^2 - \rho_2^2) + ia_3 \boldsymbol{\rho}_1 \wedge \boldsymbol{\rho}_2 \\ &= (a_1 + \gamma/2)(\rho_1^2 + \rho_2^2) - (\gamma/2)(\boldsymbol{\rho}_1 - \boldsymbol{\rho}_2)^2 \\ &\quad + ia_2(\rho_1^2 - \rho_2^2) + ia_3 \boldsymbol{\rho}_1 \wedge \boldsymbol{\rho}_2, \end{aligned}$$

$$\begin{aligned} \boldsymbol{\rho} \wedge \boldsymbol{\rho}' &= xy' - yx' \\ &= \boldsymbol{\rho} \cdot \boldsymbol{\epsilon} \boldsymbol{\rho}' \end{aligned}$$

$$\boldsymbol{\epsilon} = i\sigma_2 = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

# The Twisted Gaussian Schell Model (TGSM) beams

R. Simon and N. Mukunda

Vol. 10, No. 1/January 1993/J. Opt. Soc. Am. A

95

## Twisted Gaussian Schell-model beams

Using optically defined parameters

$$W_z(\boldsymbol{\rho}_1, \boldsymbol{\rho}_2; \nu) = \frac{I(\nu)}{2\pi\sigma_s(\nu)^2} \times \exp\left[\frac{-1}{4\sigma_s(\nu)^2}(\boldsymbol{\rho}_1^2 + \boldsymbol{\rho}_2^2) - \frac{(\boldsymbol{\rho}_1 - \boldsymbol{\rho}_2)^2}{2\sigma_g(\nu)^2}\right] \times \frac{-i}{2\chi R(\nu)}(\boldsymbol{\rho}_1^2 - \boldsymbol{\rho}_2^2) - i\frac{u(\nu)}{\chi}\boldsymbol{\rho}_1 \cdot \boldsymbol{\epsilon}\boldsymbol{\rho}_2 \quad (2.2)$$

$I$  is the intensity

$\sigma_s$  is the beam width

$\sigma_g$  is the coherence length

$R$  is radius of curvature

$u$  is the twist phase parameter

$$\chi = \lambda/2\pi$$

# The Twisted Gaussian Schell Model (TGSM) beams

## Twisted Gaussian Schell-model beams

### Effects on divergence

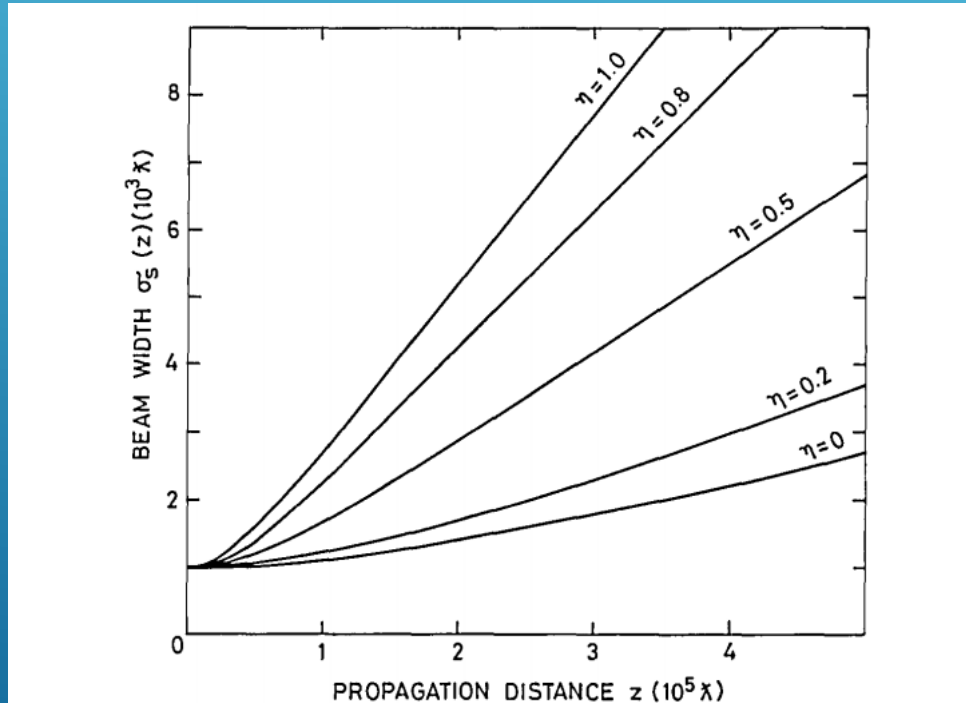


Fig. 1. Effect of the invariant twist phase on beam expansion, with  $\beta = 5$ ,  $\sigma_s(0) = 10^3 \lambda$ .

### Effect on the propagation phase

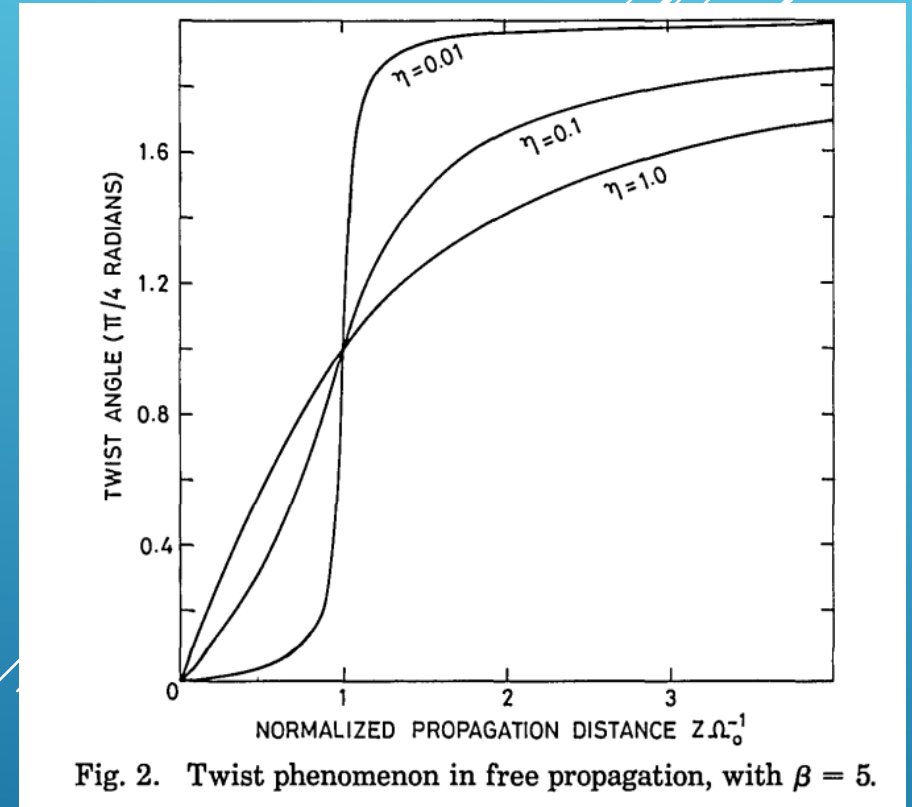


Fig. 2. Twist phenomenon in free propagation, with  $\beta = 5$ .

$$\beta = \sigma_s(z)/\sigma_g(z) = \sigma_s(0)/\sigma_g(0)$$

$\sigma_g(\nu)$  Coherence length

$\sigma_s(\nu)$  Beam width

# The Twisted Gaussian Schell Model (TGSM) beams illustrating beam rotation



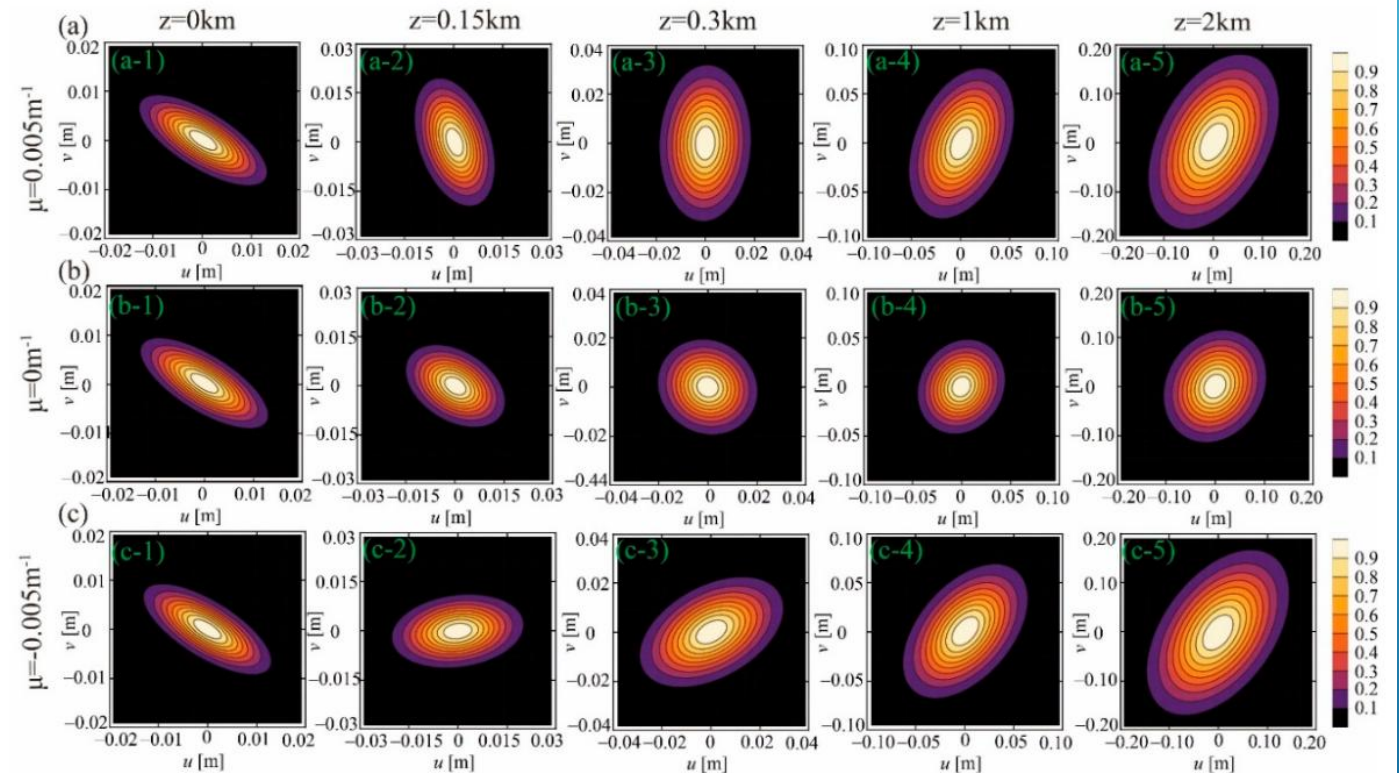
Article

## Statistical Characteristics of a Twisted Anisotropic Gaussian Schell-Model Beam in Turbulent Ocean

Yonglei Liu <sup>1,2</sup>, Yuefeng Zhao <sup>1,2</sup>, Xianlong Liu <sup>1,2</sup>, Chunhao Liang <sup>1,2</sup>, Lin Liu <sup>3</sup>, Fei Wang <sup>3</sup> and Yangjian Cai <sup>1,2,3,\*</sup>

*Photonics* 2020, 7, 37;

doi:10.3390/photonics7020037



# TGSM beams motivation

Classical Optics: Robustness against Propagation in turbulent media

F. Wang and Y. Cai, "Second-order statistics of a twisted Gaussian Schell-model beam in turbulent atmosphere," *Opt. Express*, vol. 18, p. 24661, 2010.

M. Zhou, W. Fan, and G. Wu, "Evolution properties of the orbital angular momentum spectrum of twisted Gaussian Schell-model beams in turbulent atmosphere," *J. Opt. Soc. Am. A*, vol. 37, p. 142, 2020.

Y. Liu, X. Liu, L. Liu, F. Wang, Y. Zhang, and Y. Cai, "Ghost imaging with a partially coherent beam carrying twist phase in a turbulent ocean: a numerical approach," *Appl. Sci.*, vol. 9, 2019, Art no. 3023.

Quantum Optics: Robustness against Propagation in turbulent media

**Samukelisiwe Purity Phehlukwayo , Marie Louise Umuhire , Yaseera Ismail, Stuti Joshi , and Francesco Petruccione , Influence of coincidence detection of a biphoton state through free-space atmospheric turbulence using a partially spatially coherent pump, *Phys. Rev. A* 102, 033732 (2020)**

Quantum Optics:  
Boosting quantum entanglement

L. Hutter, G. Lima, and S. P. Walborn, "Boosting entanglement generation in down-conversion with incoherent illumination," *Phys. Rev. Lett.*, vol. 125, p. 193602, 2020.

# TGSM beams: generation

1818 J. Opt. Soc. Am. A/Vol. 11, No. 6/June 1994

## Interpretation and experimental demonstration of twisted Gaussian Schell-model beams

Ari T. Friberg, Eero Tervonen, and Jari Turunen

Department of Technical Physics, Helsinki University of Technology, FIN- 02150 Espoo, Finland

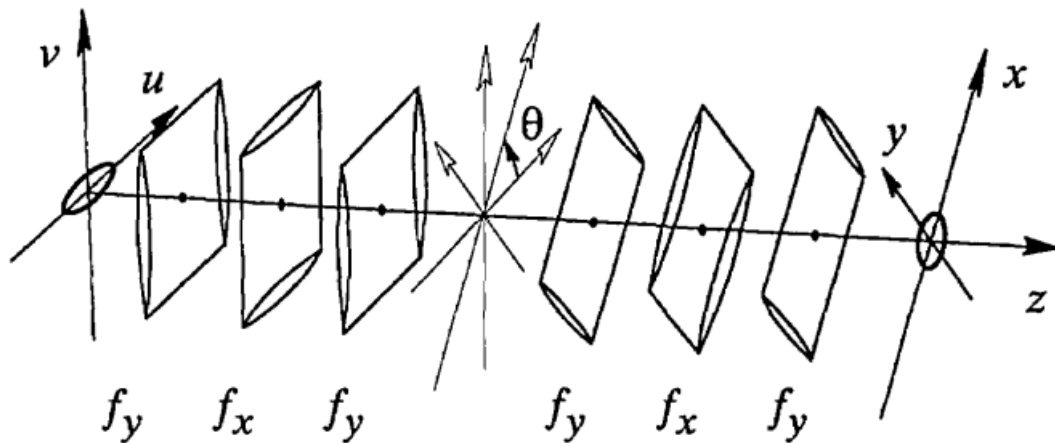


Fig. 1. Astigmatic optical lens system used for converting an anisotropic GSM beam into a twisted GSM beam.

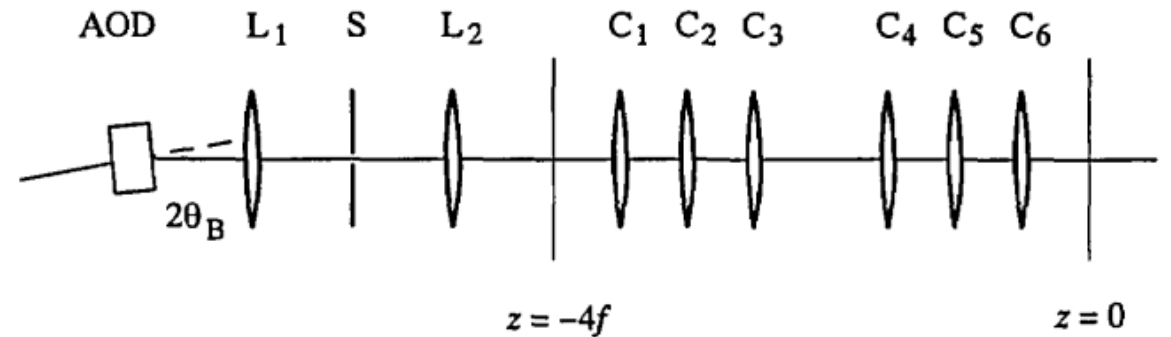
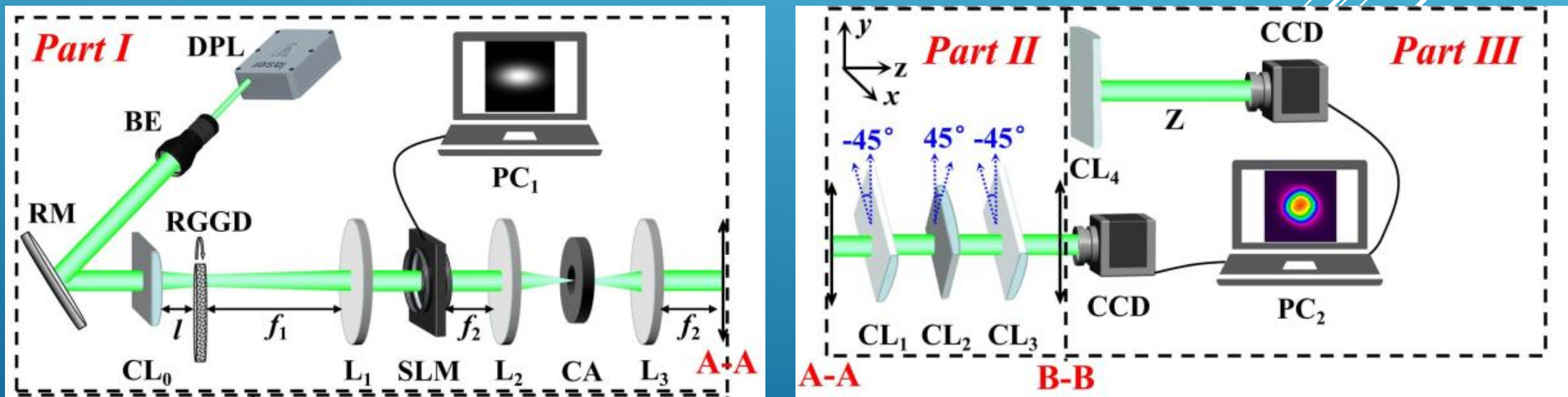


Fig. 5. Experimental arrangement: AOD, acousto-optic deflector;  $L_1$  and  $L_2$ , spherical lenses;  $C_1$ – $C_6$ , cylindrical lenses; S, spatial filter.



### Generating bona fide twisted Gaussian Schell-model beams

HAIYUN WANG,<sup>1</sup> XIAOFENG PENG,<sup>1</sup> LIN LIU,<sup>1,5</sup> FEI WANG,<sup>1,6</sup> YANGJIAN CAI,<sup>1,2,7</sup> AND SERGEY A. PONOMARENKO<sup>3,4</sup>



**Fig. 1.** Experimental setup for generating a TGSM beam. DPSS, diode-pumped solid-state laser; BE, beam expander; RM, reflecting mirror; CL<sub>0</sub>, CL<sub>1</sub>, CL<sub>2</sub>, CL<sub>3</sub>, and CL<sub>4</sub>, thin cylindrical lenses;

RGGD, rotating ground glass disk; L<sub>1</sub>, L<sub>2</sub>, and L<sub>3</sub>, thin lenses; SLM, spatial light modulator; CA, circular aperture; CCD, charge-coupled device; PC<sub>1</sub> and PC<sub>2</sub>, personal computers.

# TGSM beams: generation

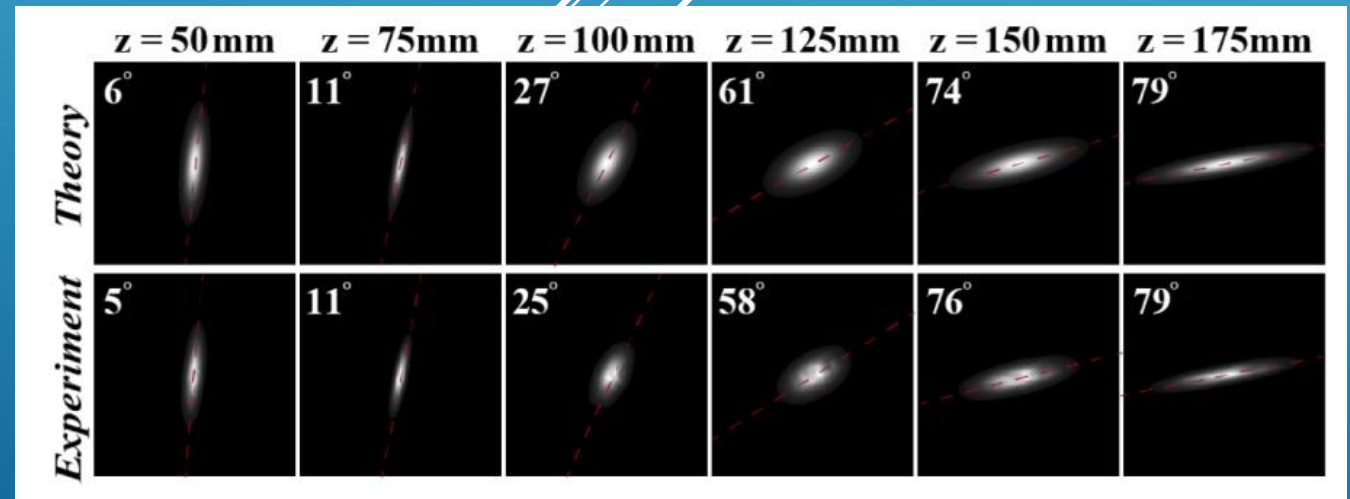
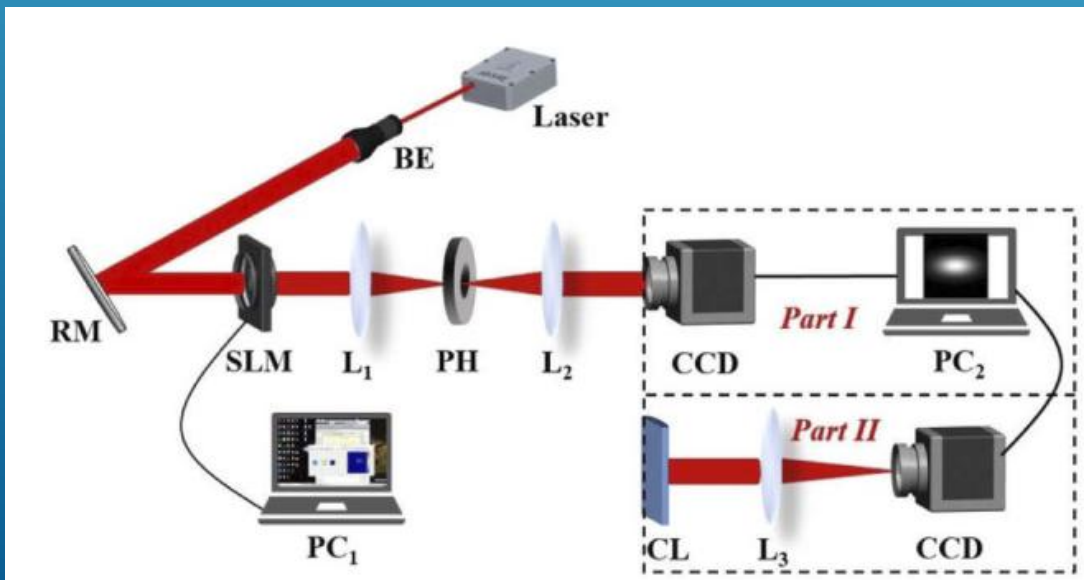
Research Article

Vol. 29, No. 25/6 Dec 2021 / *Optics Express* 41964

**Optics EXPRESS**

## Generating a twisted Gaussian Schell-model beam with a coherent-mode superposition

YUE ZHANG,<sup>1</sup> XUAN ZHANG,<sup>1</sup> HAIYUN WANG,<sup>1</sup> YAN YE,<sup>2</sup> LIN LIU,<sup>1,4</sup>  
YAHONG CHEN,<sup>1,5</sup>  FEI WANG,<sup>1,6</sup> AND YANGJIAN CAI<sup>1,3,7</sup>



# TGSM beams: generation

DE GRUYTER

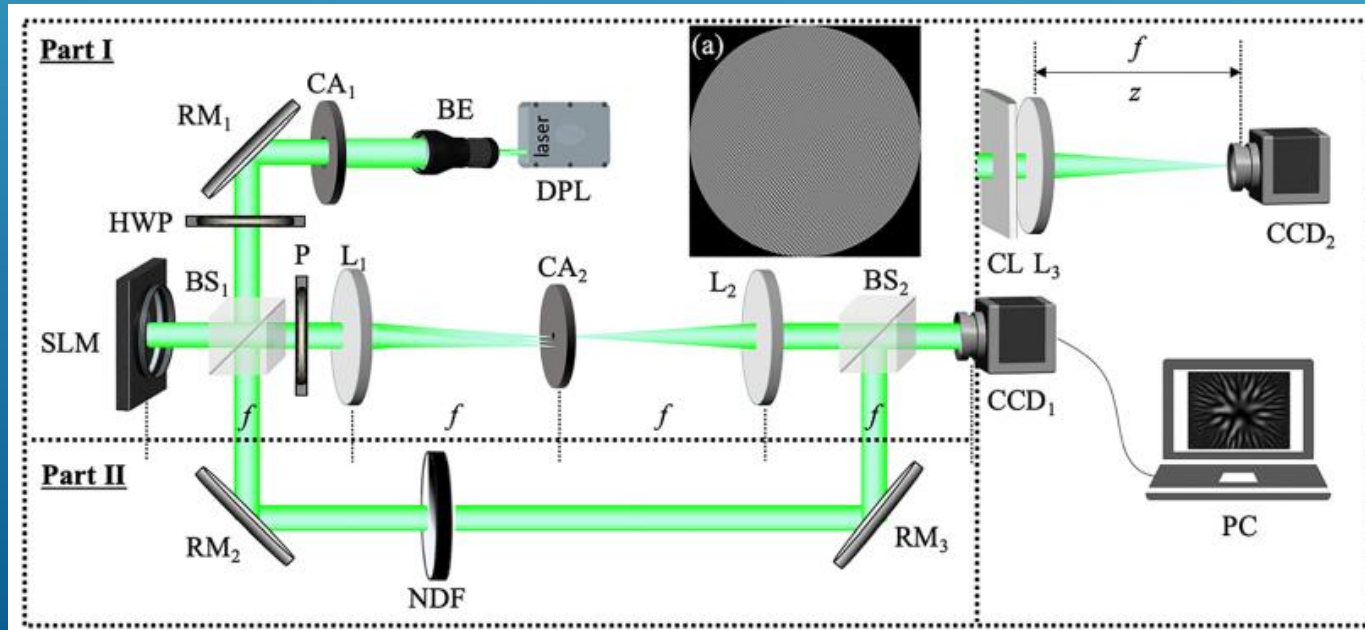
Nanophotonics 2022; 11(4): 689–696



## Research Article

Haiyun Wang, Xiaofeng Peng, Hao Zhang, Lin Liu\*, Yahong Chen\*, Fei Wang\* and Yangjian Cai\*

# Experimental synthesis of partially coherent beam with controllable twist phase and measuring its orbital angular momentum



# TGSM beams: generation

DE GRUYTER

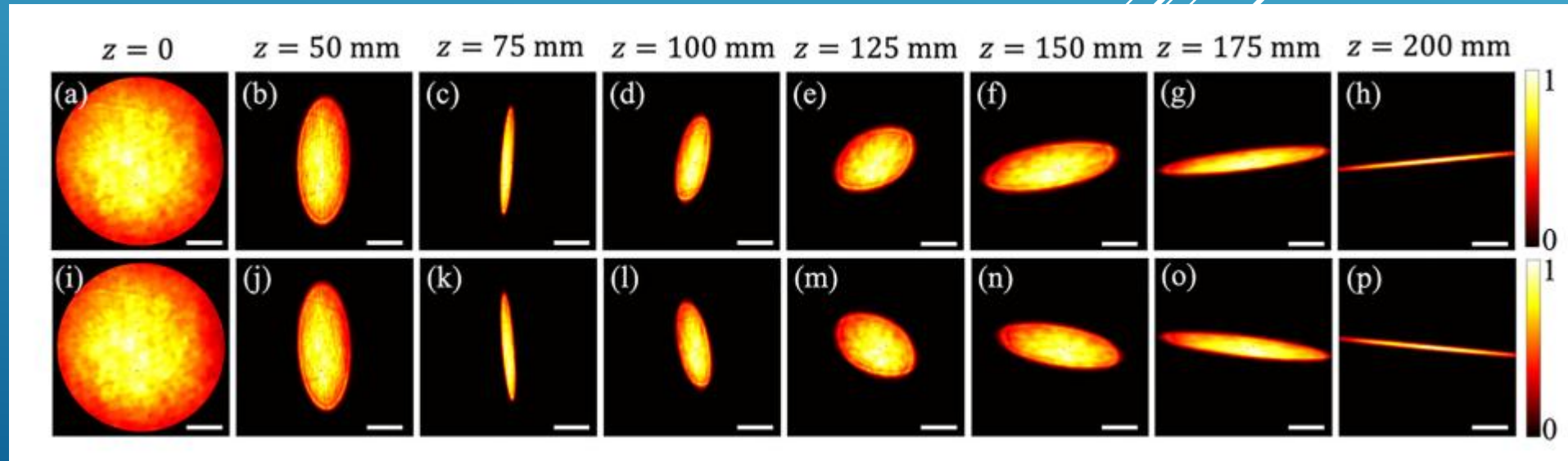
Nanophotonics 2022; 11(4): 689–696



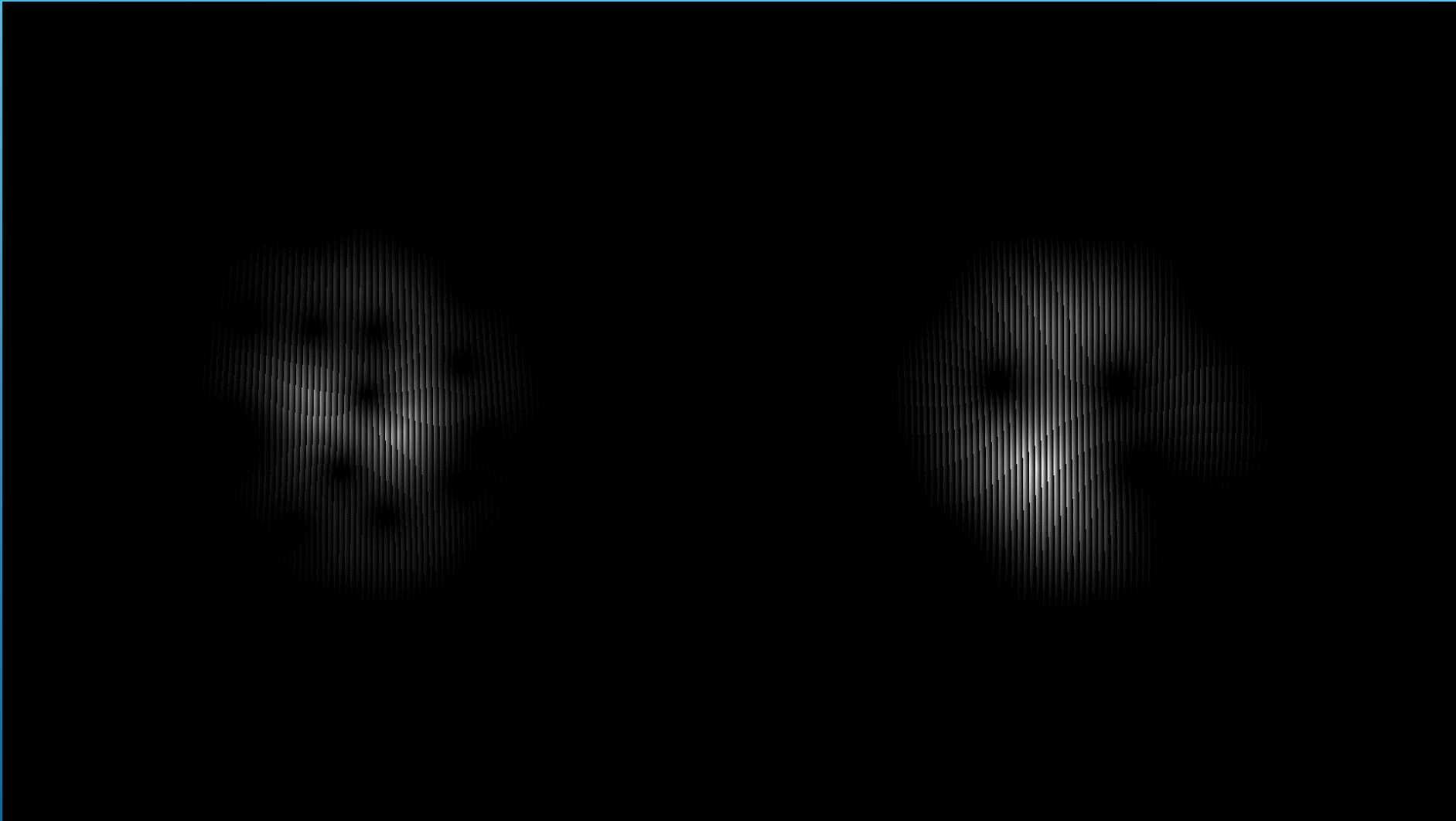
## Research Article

Haiyun Wang, Xiaofeng Peng, Hao Zhang, Lin Liu\*, Yahong Chen\*, Fei Wang\* and Yangjian Cai\*

## Experimental synthesis of partially coherent beam with controllable twist phase and measuring its orbital angular momentum

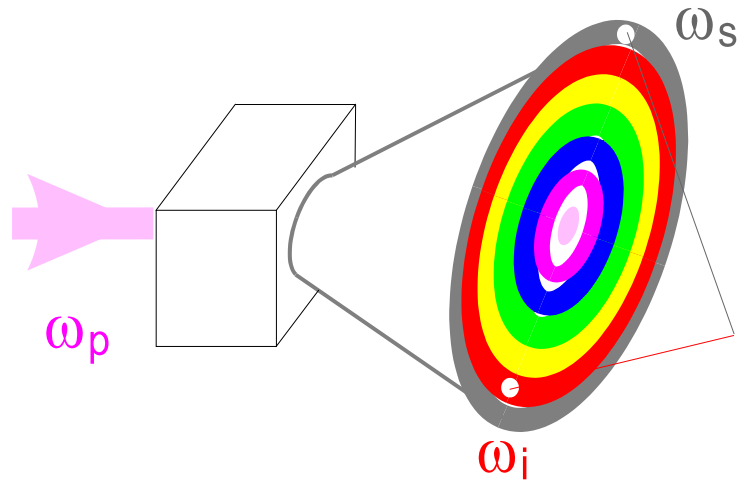


# SLM movie method for generating TGSM beams



# Parametric Down-conversion

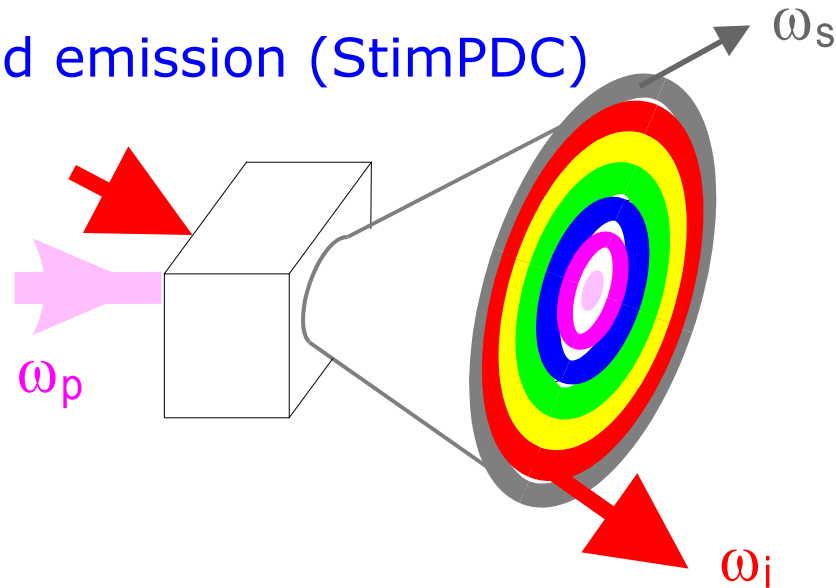
## Spontaneous emission (SPDC)



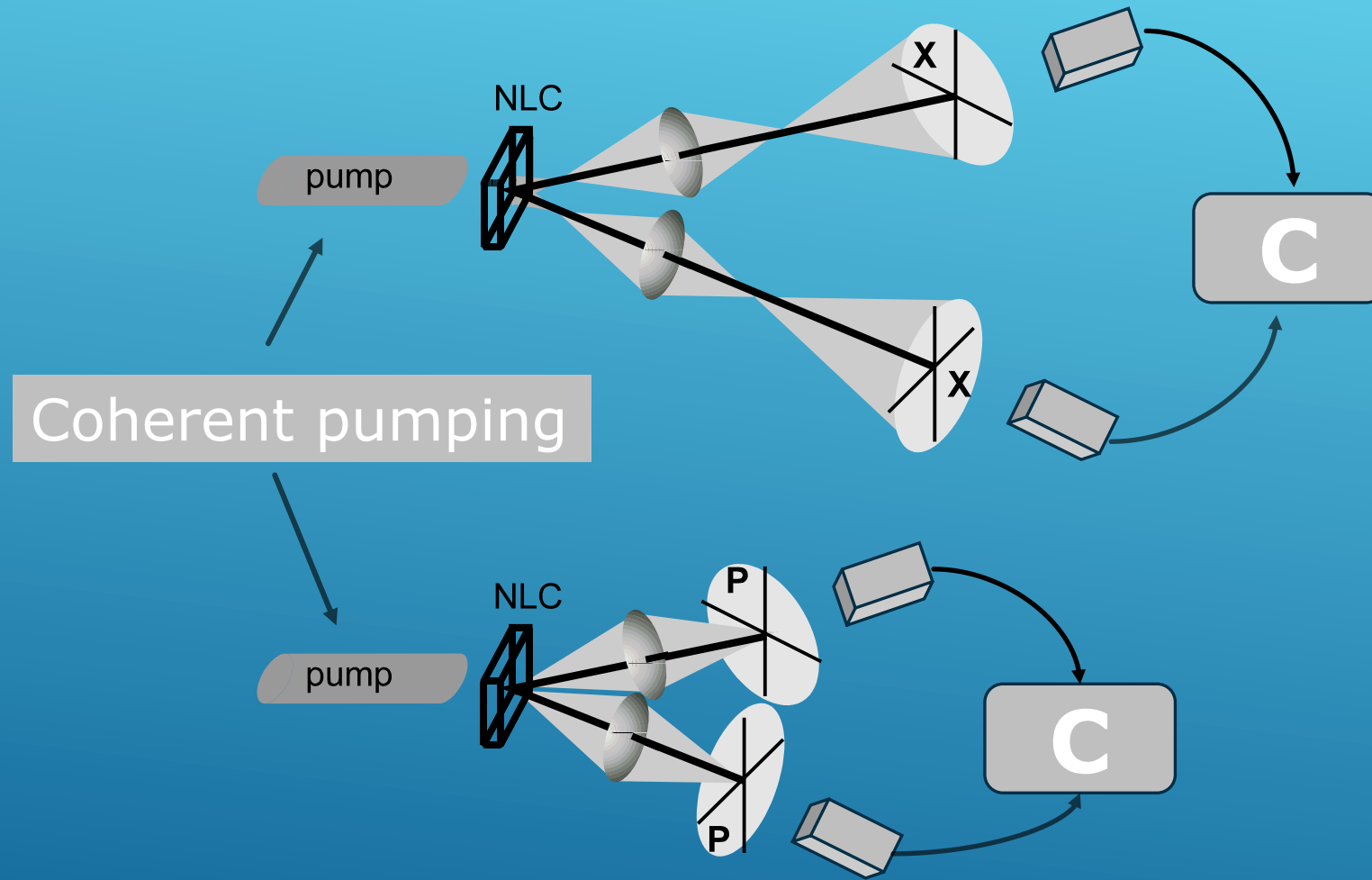
$$\hbar\omega_p = \hbar\omega_i + \hbar\omega_s$$

$$\vec{k}_p = \vec{k}_i + \vec{k}_s$$

## Stimulated emission (StimPDC)



# Spatial Entanglement Can Be Measured And Witnessed



# Spatial Entanglement Can Be Witnessed

Lu-Ming Duan, G. Giedke, J. I. Cirac, and P. Zoller  
Phys. Rev. Lett. **84**, 2722 (2000).

DGCZ criterion



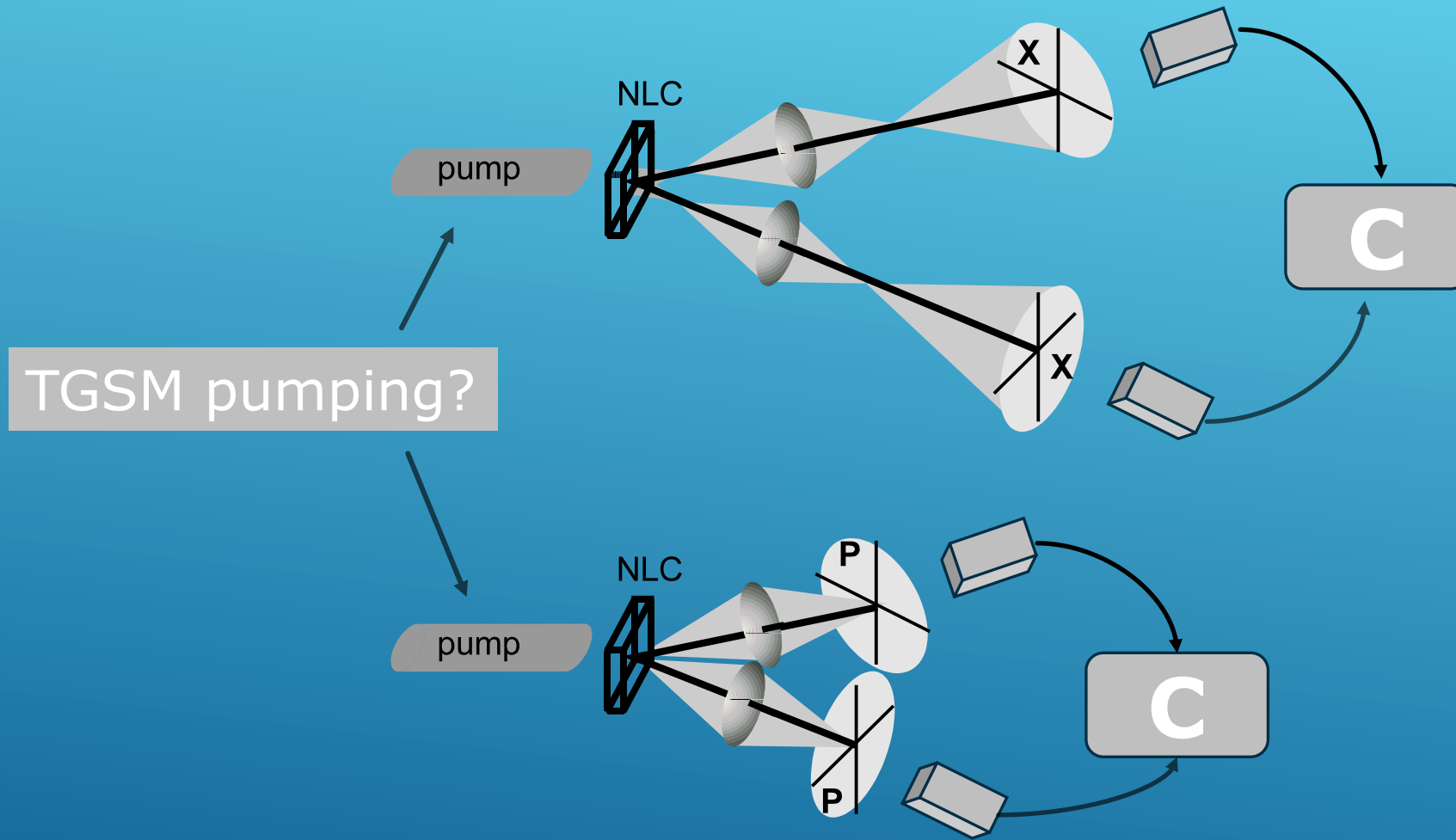
MGVT criterion

S. Mancini, V. Giovannetti, D. Vitali, and P. Tombesi  
Phys. Rev. Lett. **88**, 120401 (2002).





# Spatial Entanglement Can Be Measured And Witnessed

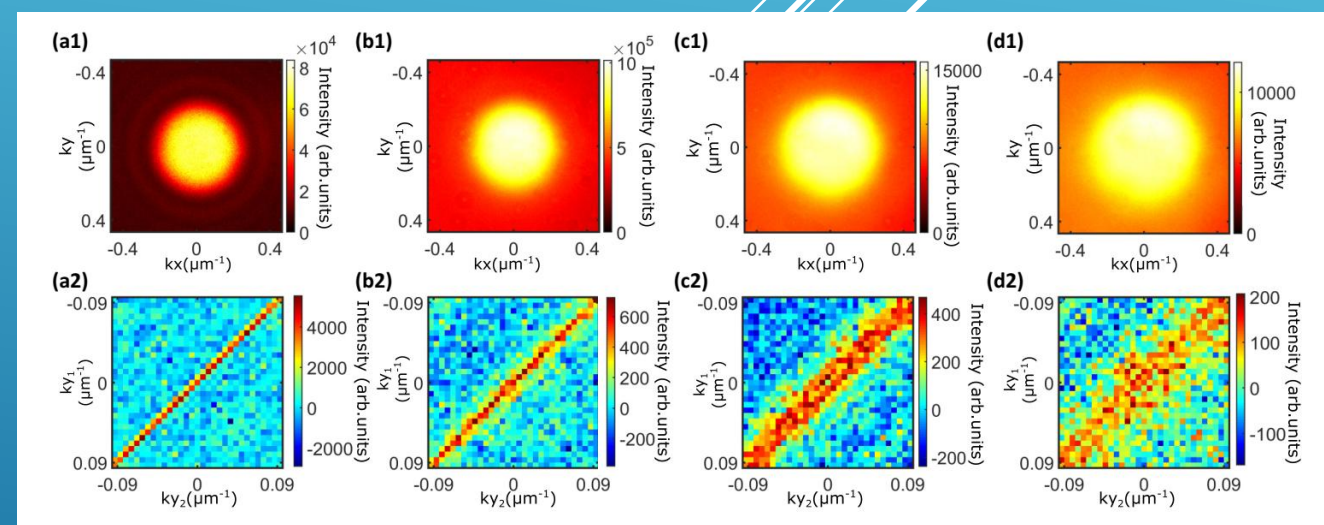
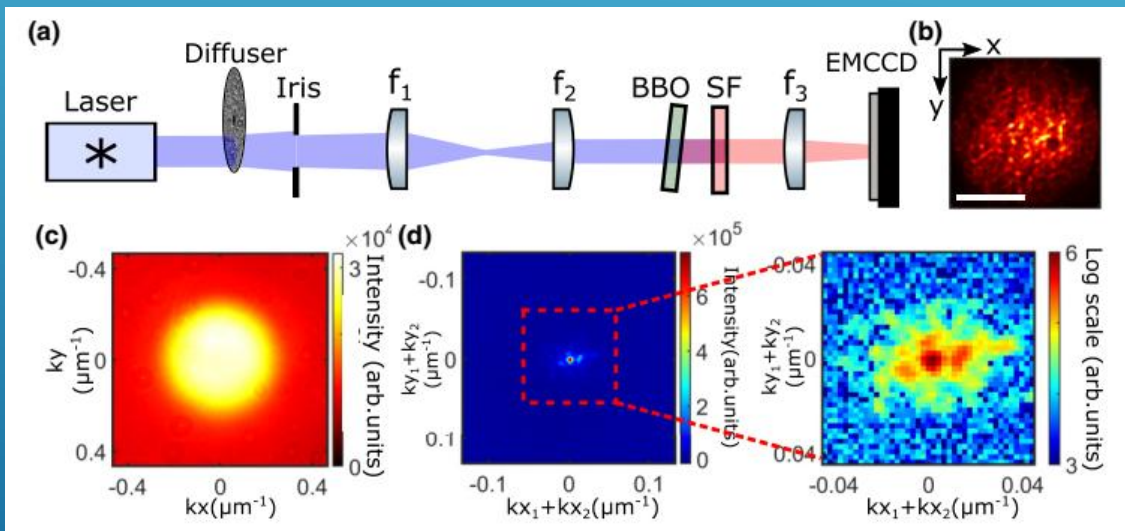


# Incoherence degrades spatial correlations

PHYSICAL REVIEW A **99**, 053831 (2019)

**Spatially entangled photon-pair generation using a partial spatially coherent pump beam**


Hugo Defienne\* and Sylvain Gigan

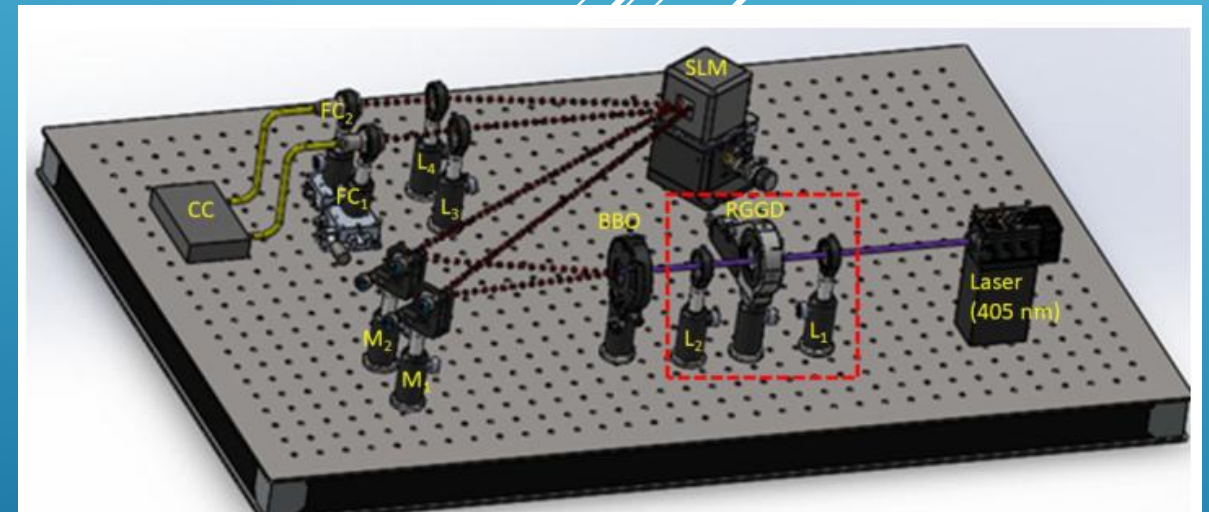
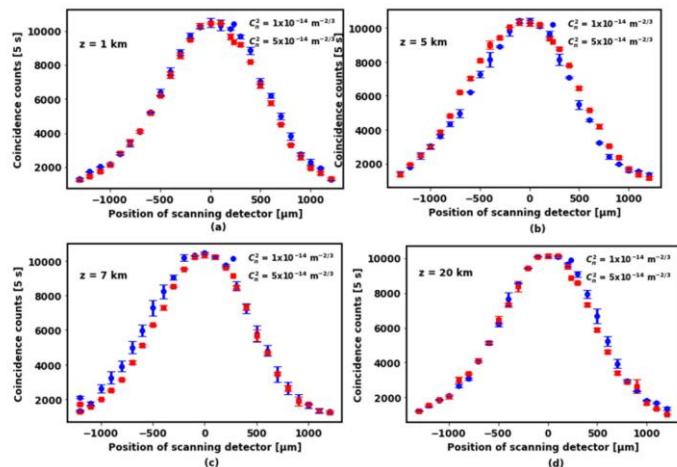
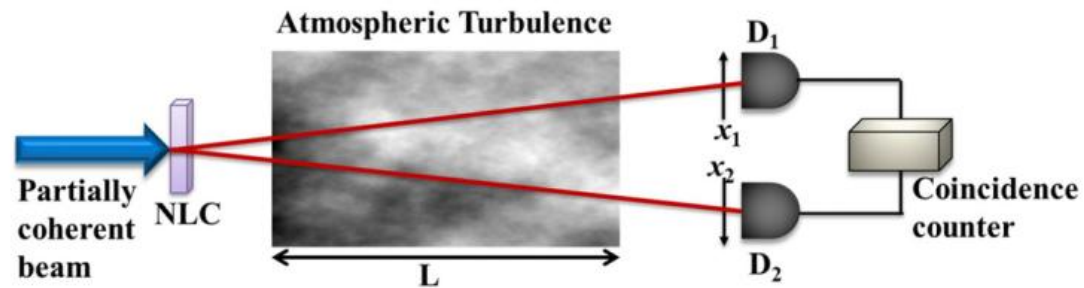


# Pumping SPDC with partially coherent pump

PHYSICAL REVIEW A **102**, 033732 (2020)

**Influence of coincidence detection of a biphoton state through free-space atmospheric turbulence using a partially spatially coherent pump**

Samukelisiwe Purity Phehlukwayo <sup>1</sup>, Marie Louise Umuhire <sup>1</sup>, Yaseera Ismail,<sup>1,\*</sup> Stuti Joshi <sup>2,†</sup>  
and Francesco Petruccione<sup>1,3,‡</sup>



# Incoherence degrades spatial correlations

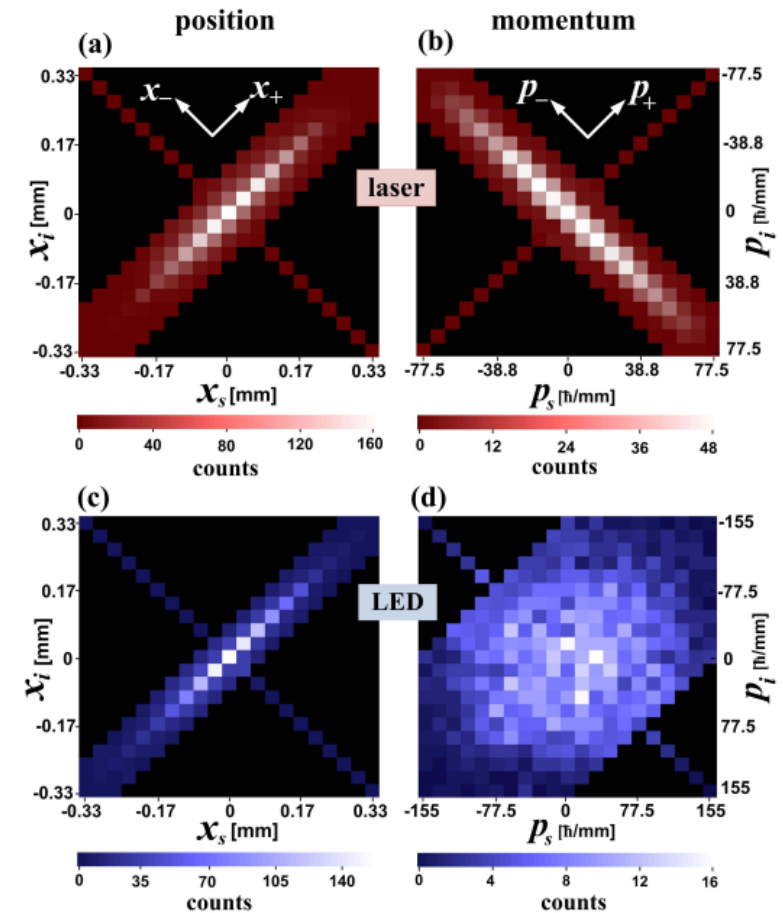
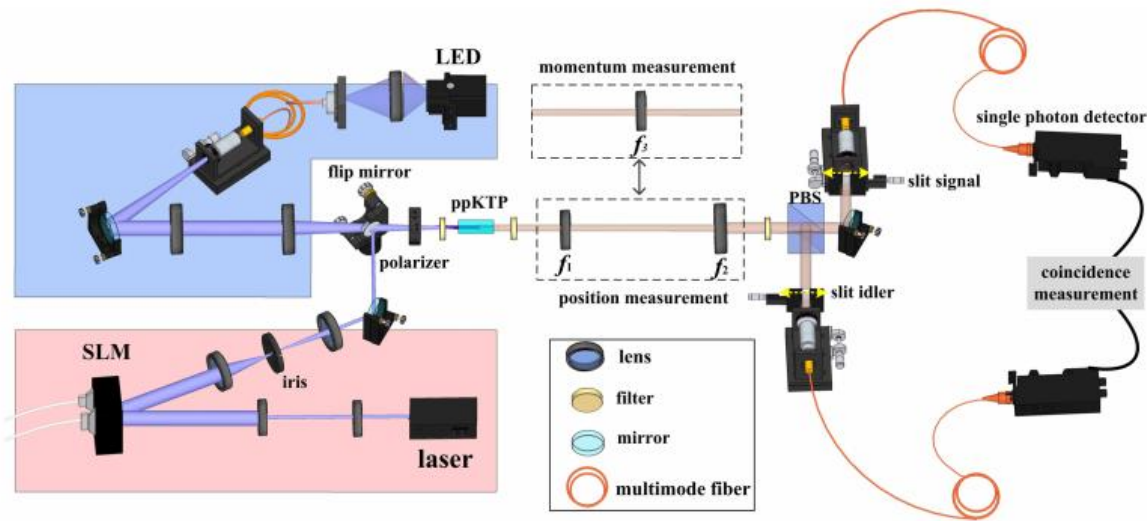
Research Article

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Optics EXPRESS

## Influence of pump coherence on the generation of position-momentum entanglement in optical parametric down-conversion

WUHONG ZHANG,<sup>1,2</sup> ROBERT FICKLER,<sup>2,3</sup> ENNO GIESE,<sup>2,4,6</sup>  
LIXIANG CHEN,<sup>1,7</sup> AND ROBERT W. BOYD<sup>2,5</sup>



# TGSM beam pumping can increase entanglement!

PHYSICAL REVIEW LETTERS **125**, 193602 (2020)

## Boosting Entanglement Generation in Down-Conversion with Incoherent Illumination

Lucas Hutter<sup>1,2</sup>, G. Lima<sup>3,4</sup> and S. P. Walborn<sup>1,3,4</sup>

A TGSM beam is therefore uniquely characterized by its CM [23,24]:

$$T = \begin{pmatrix} \sigma^2 & -\frac{k\sigma^2}{R} & 0 & ku\sigma^2 \\ -\frac{k\sigma^2}{R} & \tau^2 & -ku\sigma^2 & 0 \\ 0 & -ku\sigma^2 & \sigma^2 & -\frac{k\sigma^2}{R} \\ ku\sigma^2 & 0 & -\frac{k\sigma^2}{R} & \tau^2 \end{pmatrix}. \quad (1)$$

$\sigma$  is the beam waist

$$\tau^2 = (1/\delta^2) + (1/4\sigma^2) + k^2[(\sigma^2/R^2) + u^2\sigma^2]$$

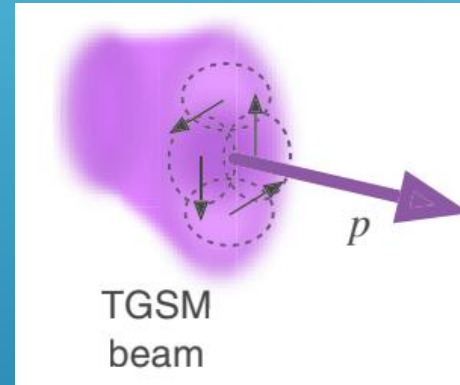
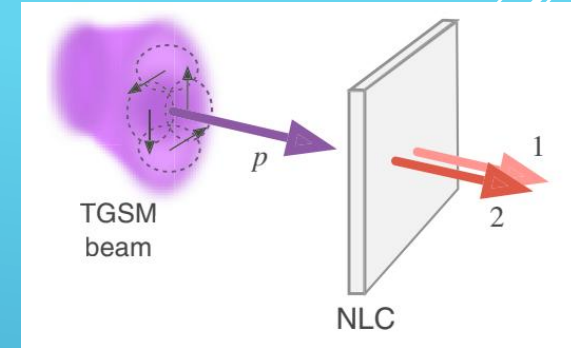
is the variance of the wave vector distribution,

$\delta$  is the transverse coherence length,

$R$  is the radius of curvature

$u$  is the so-called twist phase

$k$  is the wave number



# TGSM beam pumping can increase entanglement!

Two-photon covariance matrix

$$V_{12} = \begin{pmatrix} A & C \\ C^T & B \end{pmatrix}, \quad (4)$$

$A$  ( $B$ ) refer to photon 1 (2)     $A = B$

$$A = \begin{pmatrix} \sigma^2 + \sigma_-^2 & -\frac{k\sigma^2}{2R} & 0 & \frac{ku\sigma^2}{2} \\ -\frac{k\sigma^2}{2R} & \frac{1}{4}(\tau^2 + \Delta_-^2) & -\frac{ku\sigma^2}{2} & 0 \\ 0 & -\frac{ku\sigma^2}{2} & \sigma^2 + \sigma_-^2 & -\frac{k\sigma^2}{2R} \\ \frac{ku\sigma^2}{2} & 0 & -\frac{k\sigma^2}{2R} & \frac{1}{4}(\tau^2 + \Delta_-^2) \end{pmatrix}$$

$$\sigma_-^2 = 9L/10k$$

$L$  is the length of the crystal

$$\Delta_-^2 = 3k/2L$$

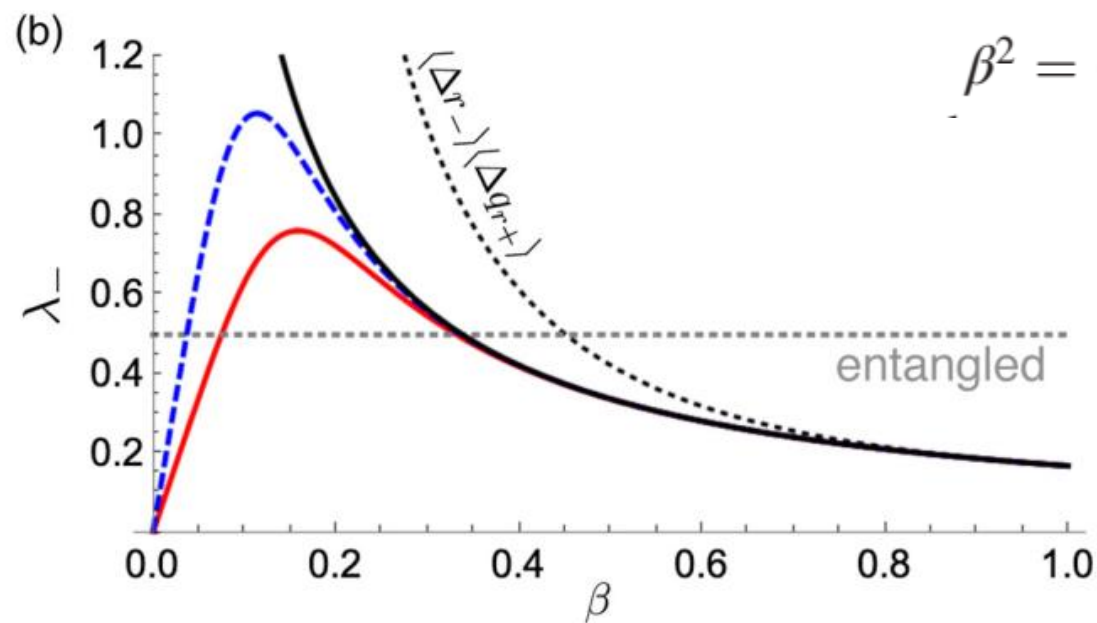
$$C = \begin{pmatrix} \sigma^2 - \sigma_-^2 & -\frac{k\sigma^2}{2R} & 0 & \frac{ku\sigma^2}{2} \\ -\frac{k\sigma^2}{2R} & \frac{1}{4}(\tau^2 - \Delta_-^2) & -\frac{ku\sigma^2}{2} & 0 \\ 0 & -\frac{ku\sigma^2}{2} & \sigma^2 - \sigma_-^2 & -\frac{k\sigma^2}{2R} \\ \frac{ku\sigma^2}{2} & 0 & -\frac{k\sigma^2}{2R} & \frac{1}{4}(\tau^2 - \Delta_-^2) \end{pmatrix}$$

The symplectic eigenvalues of Eq. (4) are twofold degenerate and given by [37]

$$\lambda_{\pm} = \frac{1}{\sqrt{2}} \left| \sqrt{a_{+} \pm \sqrt{4k^2 \Delta_{-}^2 \sigma_{-}^2 \sigma^4 \left( u^2 + \frac{1}{R^2} \right) + a_{-}^2}} \right|, \quad (8)$$

where  $a_{\pm} = \tau^2 \sigma_{-}^2 \pm \Delta_{-}^2 \sigma^2$ .

Entanglement is confirmed when  $\lambda_{-} < 1/2$  (gray horizontal plane). The SPDC parameters are  $R = \infty$ ,  $\lambda_p = 400$  nm,  $\sigma_p = 50$   $\mu$ m, and  $L = 1$  cm. (b) Profile plots of  $\lambda_{-}$  for normalized twist phase  $|u|/k\delta^2$  equal to zero (black solid line), 1 (red solid line), and 1/2 (blue dashed line). The dotted black curve is the near-field and far-field entanglement criteria (7).



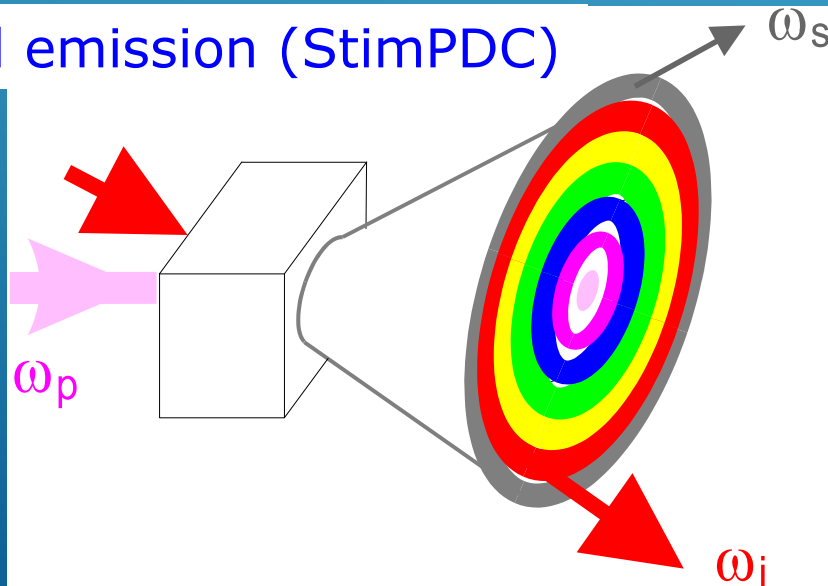
$$\beta^2 = (1 + 4\sigma^2/\delta^2)^{-1}$$

## Research Article

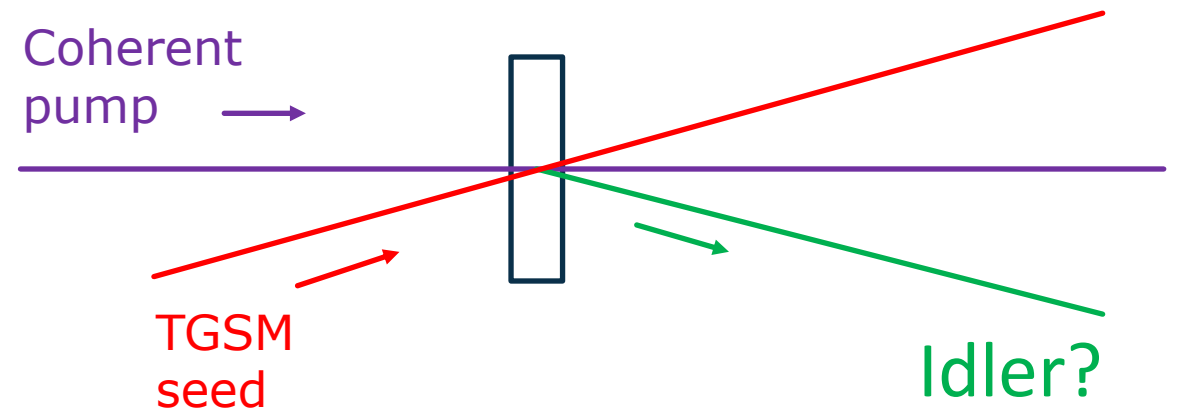
Gustavo H. dos Santos, Andre G. de Oliveira, Nara Rubiano da Silva, Gustavo Cañas, Esteban S. Gómez, Stuti Joshi, Yaseera Ismail, Paulo H. Souto Ribeiro and Stephen Patrick Walborn\*

# Phase conjugation of twisted Gaussian Schell model beams in stimulated down-conversion

Stimulated emission (StimPDC)



Coherent pump

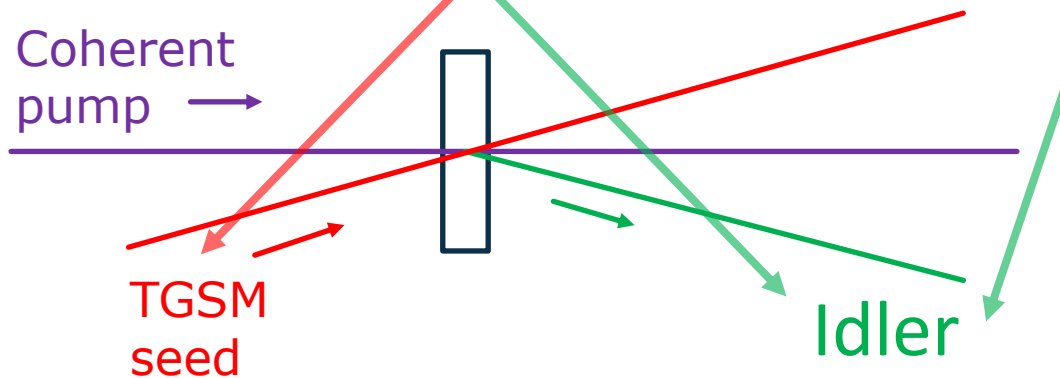




## Partially coherent StimPDC

$$W_i(\mathbf{r}, \mathbf{r}') = W_p(\mathbf{r}, \mathbf{r}')W_s^*(\mathbf{r}, \mathbf{r}').$$

Coherent pump →



## StimPDC with twisted Schell-model beams

$$W_i(\mathbf{r}, \mathbf{r}')$$

$$= A e^{-\frac{r^2+r'^2}{4w_i^2}} e^{-\frac{(r-r')^2}{2\delta_i^2}} e^{-ik_i \frac{(r-r')^2}{2R_i}} e^{-ik_i \mu_i (xy' - yx')},$$

$$w_i^2 = \frac{w_s^2 w_p^2}{w_s^2 + w_p^2},$$

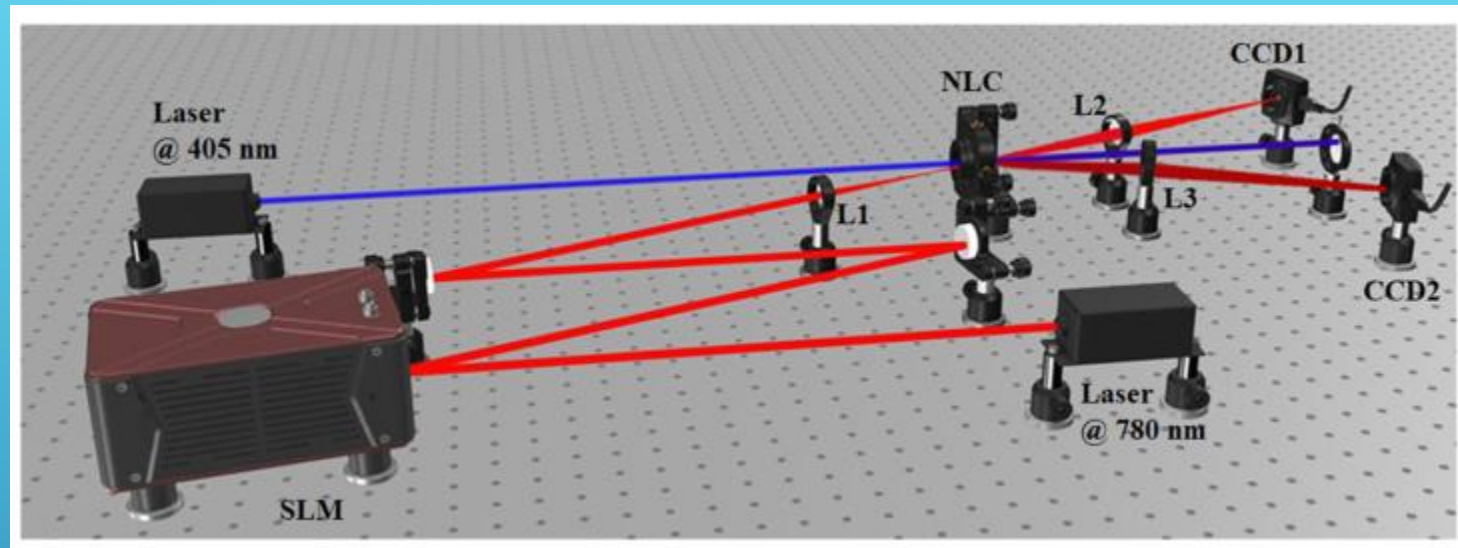
$$\delta_i^2 = \frac{\delta_s^2 \delta_p^2}{\delta_s^2 + \delta_p^2},$$

$$\frac{k_i}{R_i} = \frac{k_p}{R_p} - \frac{k_s}{R_s},$$

$$k_i \mu_i = k_p \mu_p - k_s \mu_s,$$

# Experiment

Pump @ 405 nm  
Seed @ 780 nm  
Idler @ 840 nm

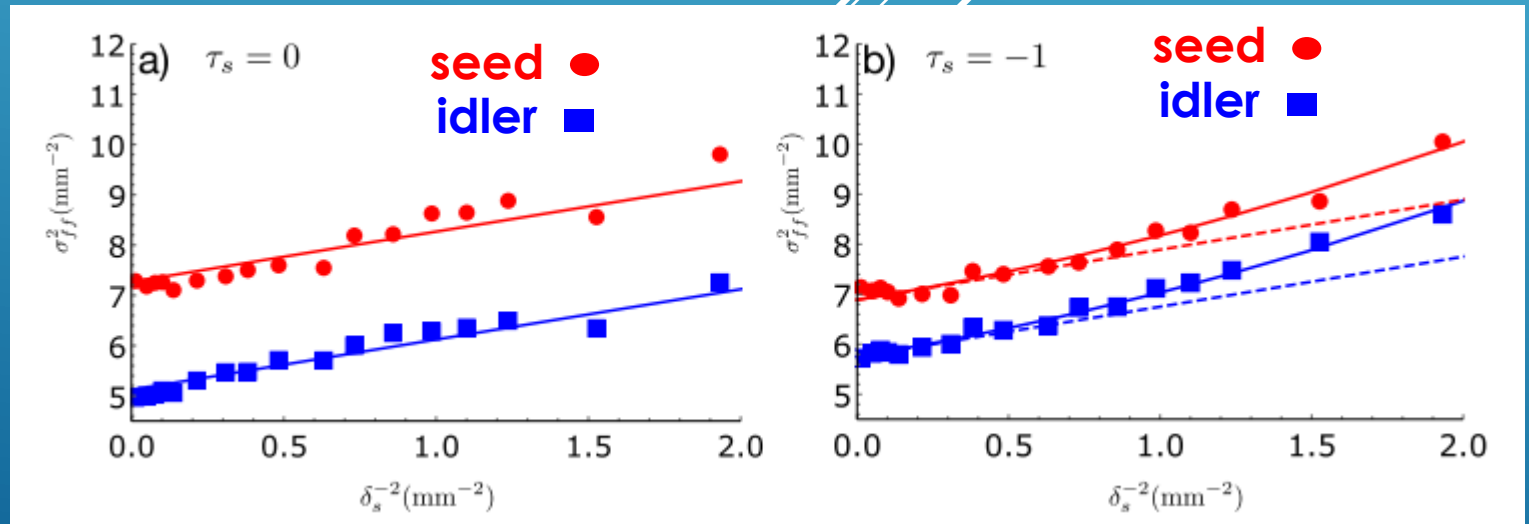


# Results

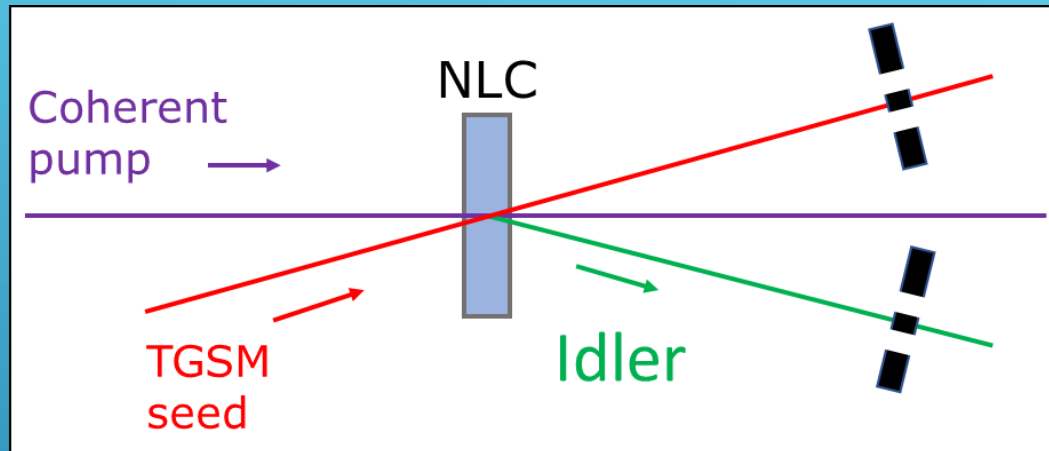
Far field variance x degree of coherence

Far field variance

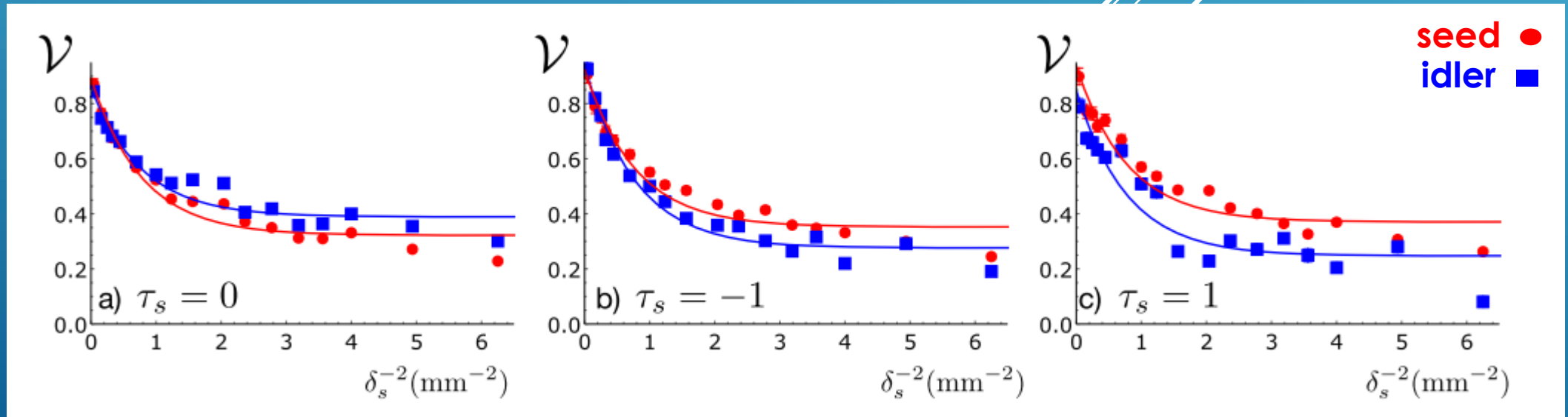
$$\sigma_{ff}^2 = \left( \frac{1}{4w^2} + \frac{k^2w^2}{R^2} \right) + \frac{1}{\delta^2} + \frac{\tau^2w^2}{\delta^4},$$



# Visibility in a double-slit interference



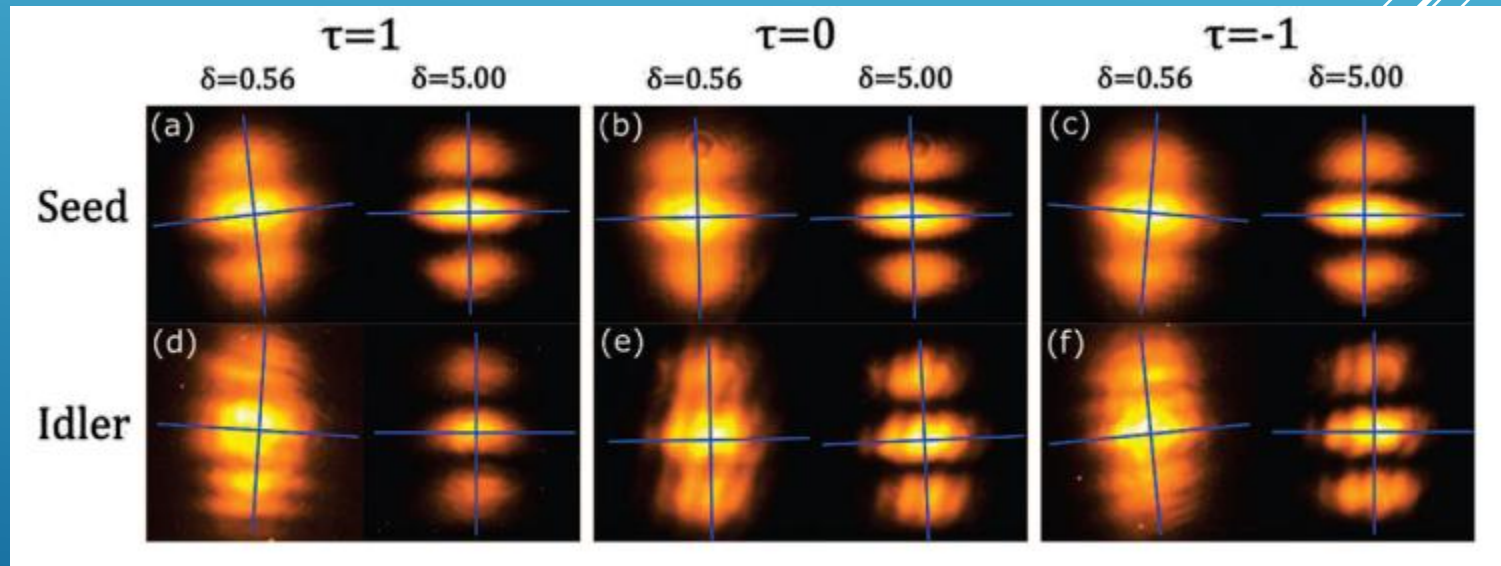
Visibility x degree of coherence



## Transfer and conjugation of twist phase

$$I(\mathbf{r}) \approx \left[ 1 + e^{-\frac{2d^2}{\delta^2}} \cos \left\{ 2dk \left( \frac{y}{f} - \mu x \right) \right\} \right]$$

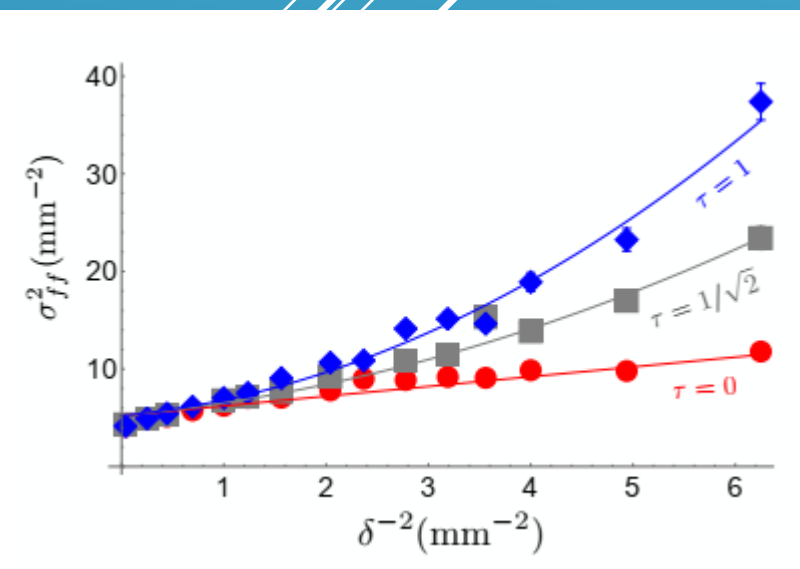
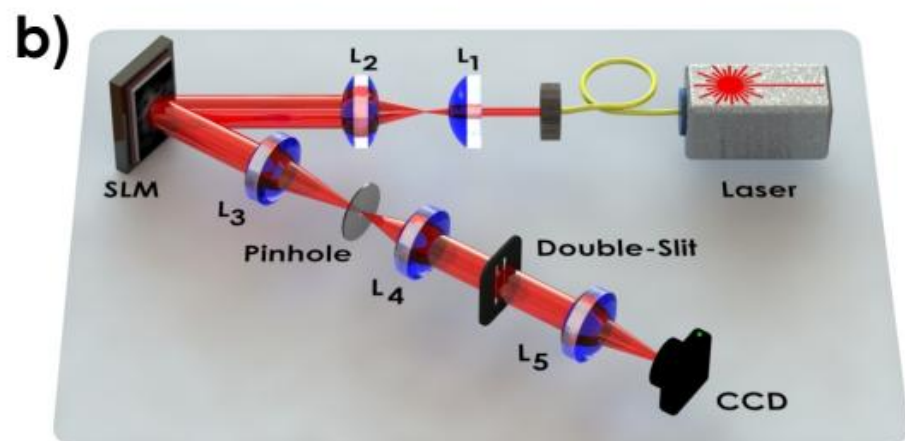
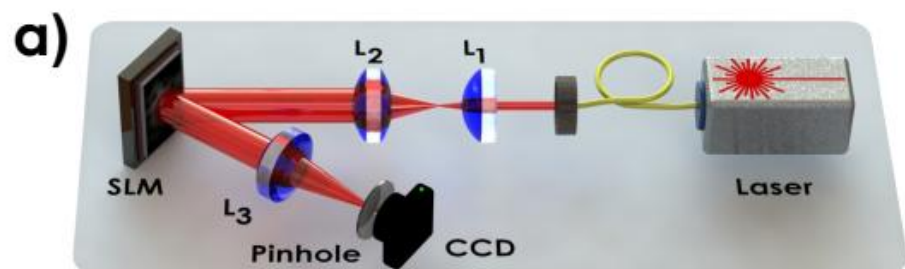
Twist in the fringe pattern



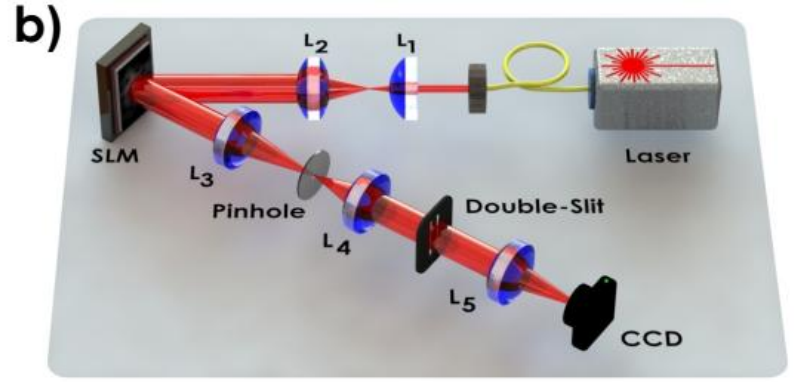
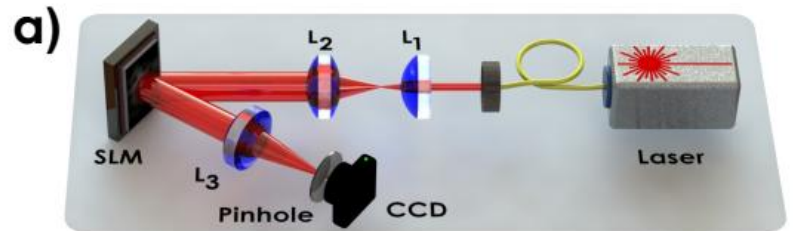
# Evaluation of twisted Gaussian Schell model beams produced with phase randomized coherent fields

G Cañas<sup>1,2</sup>, E S Gómez<sup>2,3</sup>, G H dos Santos<sup>4</sup>, A G de Oliveira<sup>4</sup>, N Rubiano da Silva<sup>4</sup>, Stuti Joshi<sup>5</sup>, Yaseera Ismail<sup>6</sup>, P H S Ribeiro<sup>4,\*</sup> and S P Walborn<sup>2,3,\*</sup>

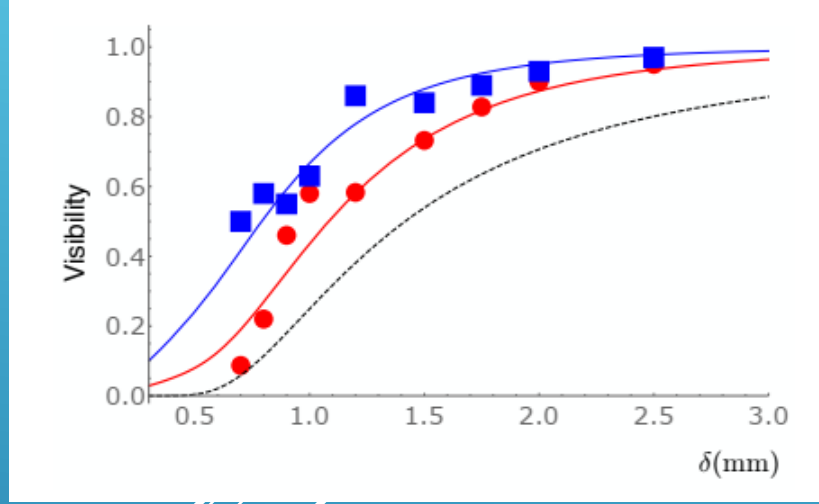
Far field mean variance versus coherence length



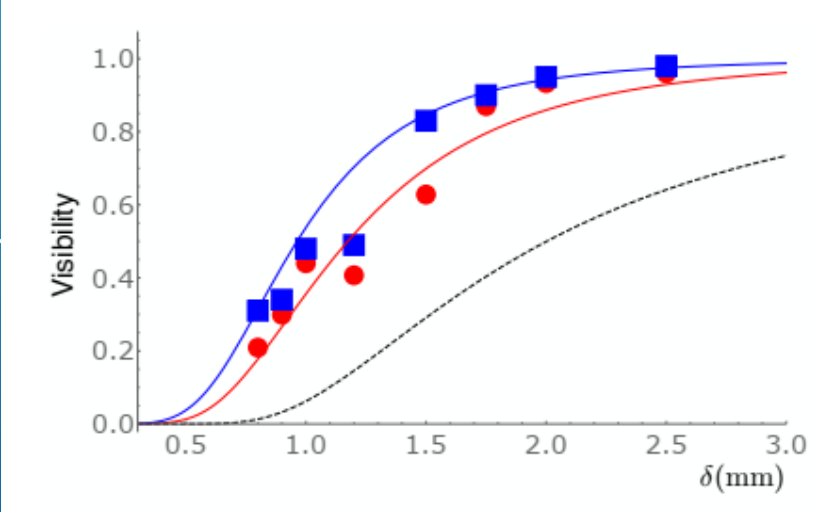
# Residual coherence in SLM movie methods



$\tau = 0$



$\tau = 1$



## Partial Coherence and Coherence Length in Stimulated Parametric Down-Conversion

G.H. dos Santos<sup>1,\*</sup>, R.C. Souza Pimenta<sup>1</sup>, R.M. Gomes<sup>2</sup>, S.P. Walborn<sup>3,4,†</sup> and P.H. Souto Ribeiro<sup>1,‡</sup>

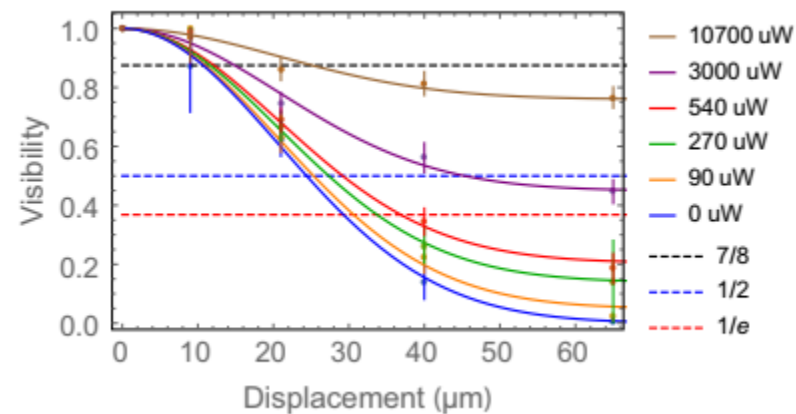
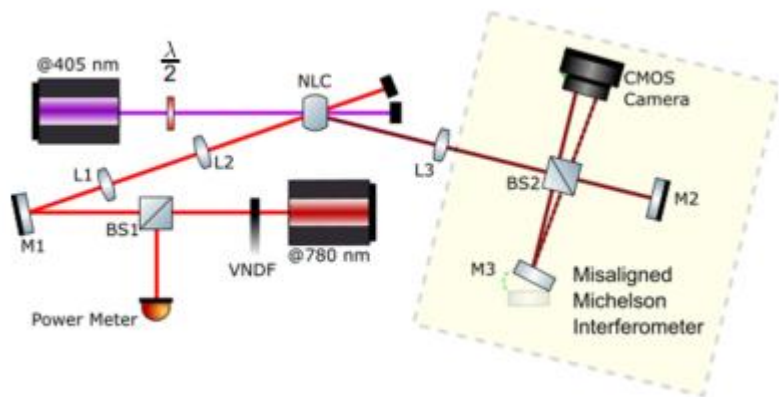


TABLE I. Normalized intensity of the stimulated component and  $\epsilon$ -coherence lengths obtained for different seed beam intensities.

$I_{\text{seed}} (\mu\text{W})$	$\beta$ (%)	$\epsilon = 1/e$ ( $\mu\text{m}$ )	$\epsilon = 1/2$ ( $\mu\text{m}$ )	$\epsilon = 7/8$ ( $\mu\text{m}$ )
0	0	29.30	24.40	10.71
90	$4.9 \pm 3.5$	30.64	25.32	11.00
270	$13.9 \pm 2.7$	33.74	27.32	11.61
540	$20.5 \pm 2.5$	36.88	29.17	12.12
3000	$44.9 \pm 1.8$	$\infty$	45.30	14.87
10 700	$76.0 \pm 1.7$	$\infty$	$\infty$	25.12

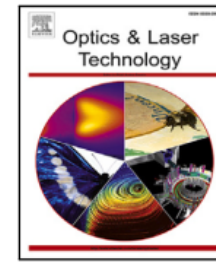


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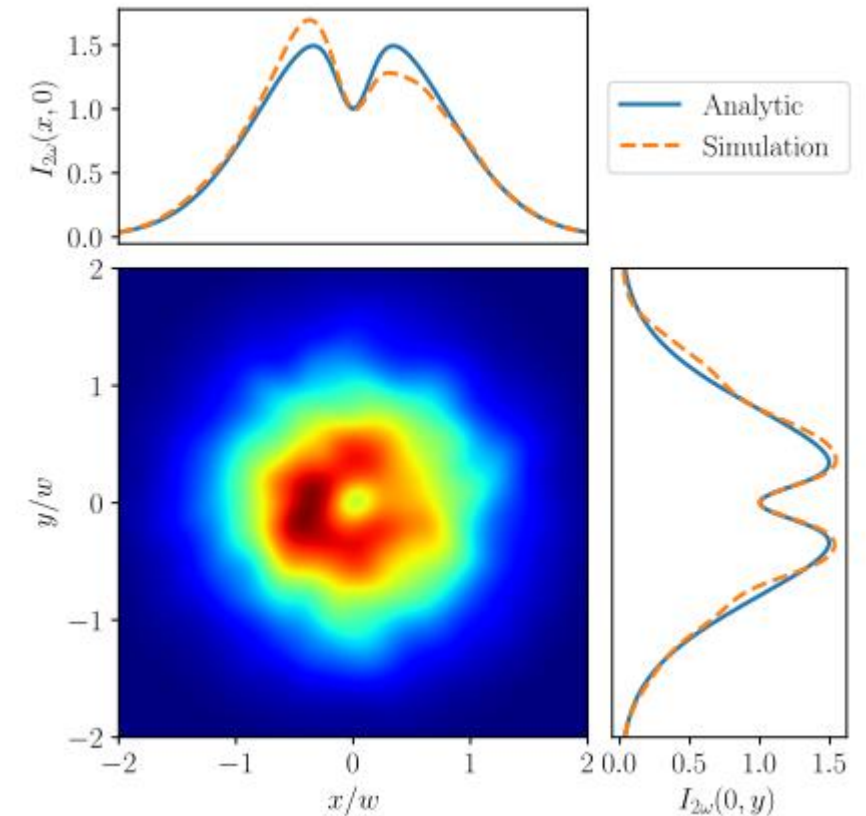
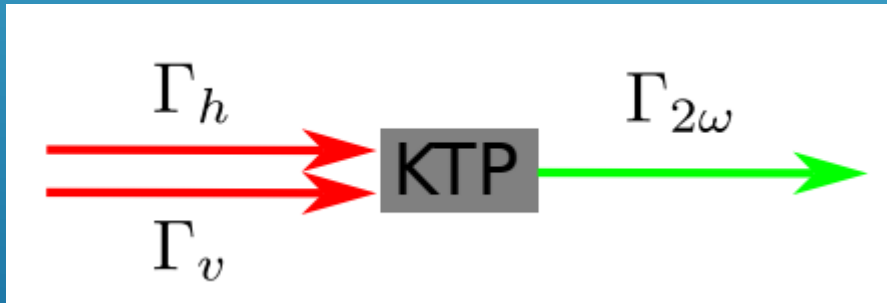
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Full length article

## Anomalous second harmonic generation of twisted Gaussian Schell model beams

M. Gil de Oliveira <sup>a,\*</sup>, A.L.S. Santos Junior <sup>a</sup>, A.C. Barbosa <sup>a</sup>, B. Pinheiro da Silva <sup>a</sup>, G.H. dos Santos <sup>b</sup>, G. Cañas <sup>c,d</sup>, P.H. Souto Ribeiro <sup>b</sup>, S.P. Walborn <sup>c,e</sup>, A.Z. Khoury <sup>a</sup>





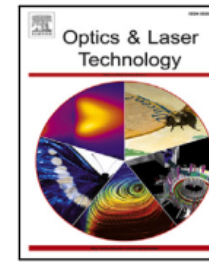


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## Anomalous second harmonic generation of twisted Gaussian Schell model beams

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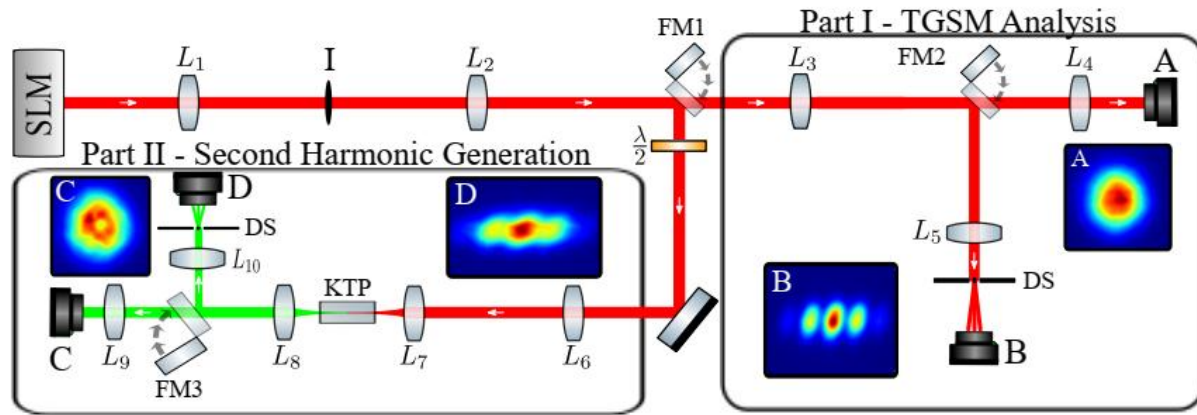
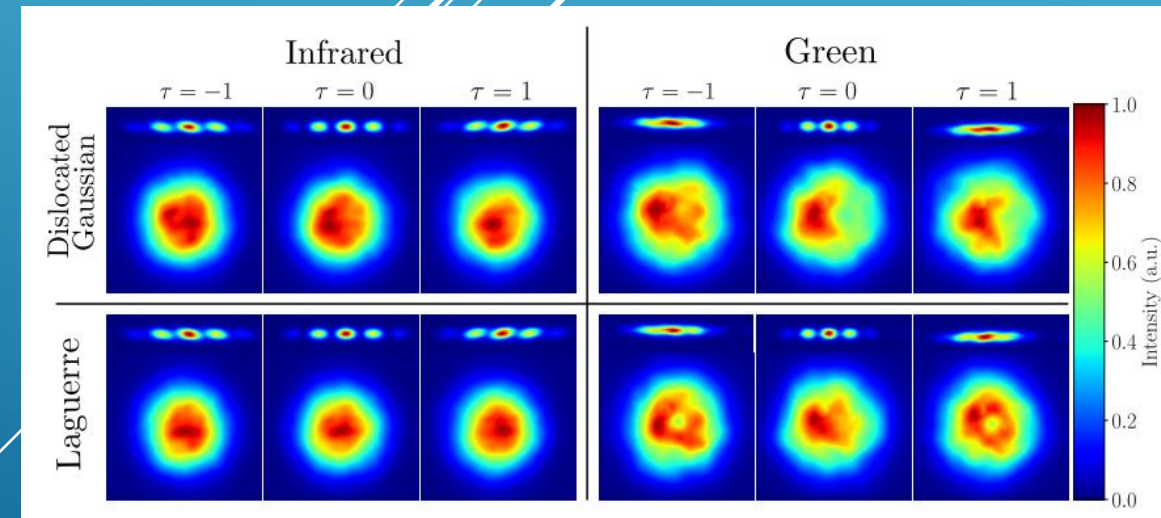


Figure 4: Experimental setup.



# Conclusions

Partial coherence GSM and TGSM beams

TGSM beams: applications to communication through turbulent medium

TGSM and SPDC applications to quantum systems

TGSM and wavelength conversion in StimPDC

StimPDC as a design for SPDC experiments in quantum regime

## Perspectives

Testing twist conservation in nonlinear parametric interactions

Testing experimentally the use of TGSM to boost the entanglement

Using TGSM beams in optical communication  
REDE RIO QUÂNTICA (A. Z. Khoury)

# Floripa

Thank you!

